

From Vorob'ev's theorem to monogamy of non-locality and local macroscopic averages

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17th May 2021

This talk is based on:

'On monogamy of non-locality and macroscopic averages'
RSB, Proceedings QPL 2014, arXiv:1412.8541 [quant-ph].

'Contextuality in quantum mechanics and beyond'
RSB, DPhil thesis, University of Oxford, 2015.

Monogamy and average macroscopic locality

- ▶ **Average macroscopic correlations from microscopic models are local**

For multipartite quantum models:

'Local realism of macroscopic correlations'

Ramanathan, Paterek, Kay, Kurzyński, Kaszlikowski, Phys. Rev. Lett. 107, 060405, 2011.

- ▶ **Monogamy of violation of Bell inequalities**

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- ▶ Connect and generalise these two results,
- ▶ providing a structural explanation related to Vorob'ev's theorem.
- ▶ We will mainly consider a simple illustrative example.

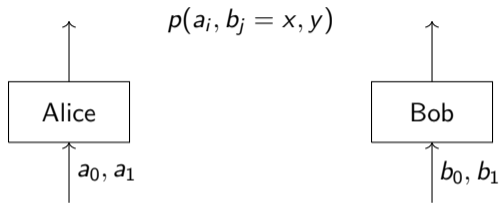
Related

Related:

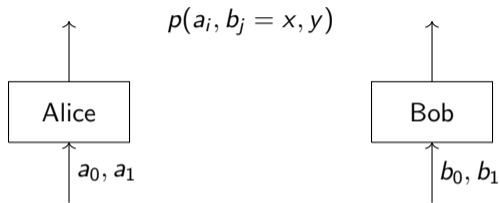
'Generalized monogamy of contextual inequalities from the no-disturbance principle'
Ramanathan, Soeda, Kurzyński, Kaszlikowski, Phys. Rev. Lett. 109, 050404, 2012.

Monogamy of non-locality

Non-locality

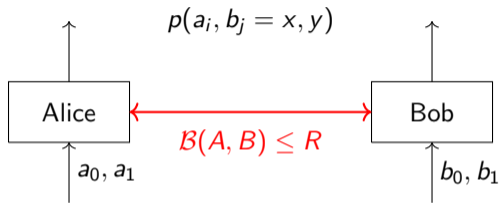


Non-locality



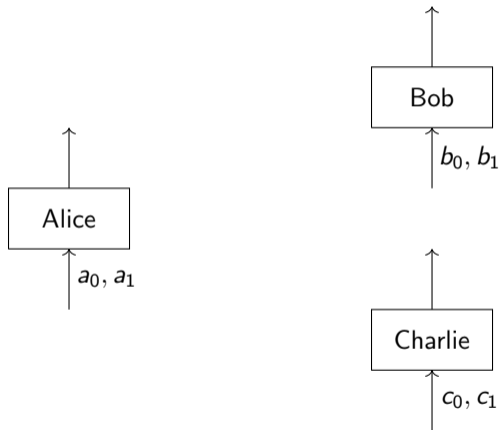
	00	01	10	11
$a_0 b_0$	$1/2$	0	0	$1/2$
$a_0 b_1$	$3/8$	$1/8$	$1/8$	$3/8$
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Monogamy of non-locality



Monogamy of non-locality

- ▶ Empirical model: no-signalling probabilities

$$p(a_i, b_j, c_k = x, y, z)$$

where x, y, z are possible outcomes.

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- ▶ Consider the subsystem composed of A and B only, given by marginalisation (in QM, partial trace):

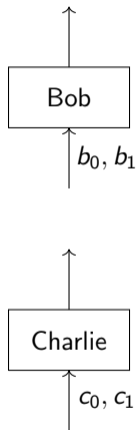
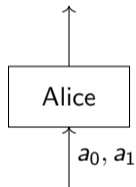
$$p(a_i, b_j = x, y) = \sum_z p(a_i, b_j, c_k = x, y, z)$$

This is independent of c_k due to no-signalling.

- ▶ Similarly define $p(a_i, c_k = x, z)$. (subsystem consisting of A and C)

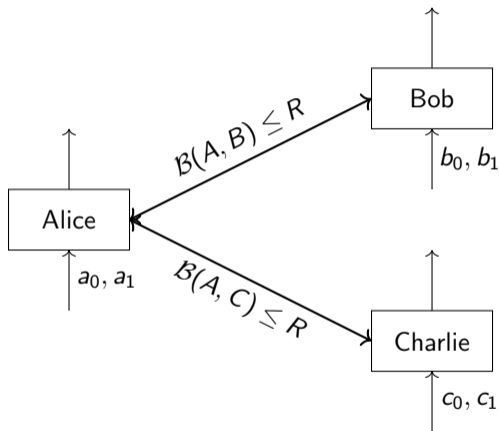
Monogamy of non-locality

Given a Bell inequality $\mathcal{B}(-, -,) \leq R$,



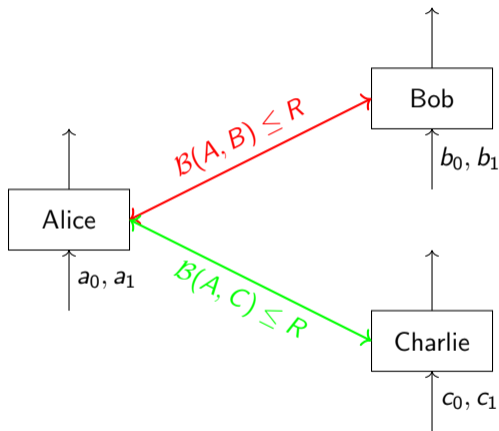
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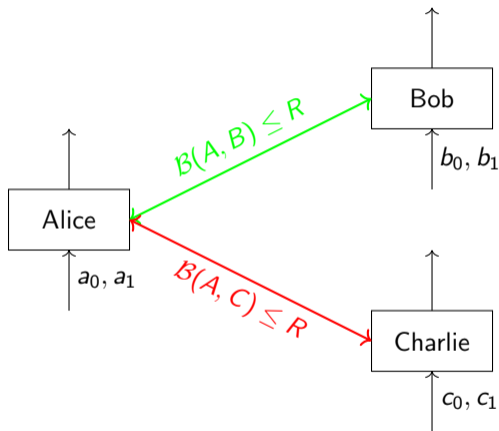
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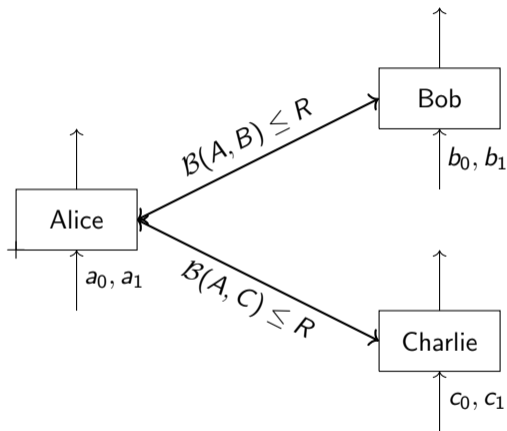
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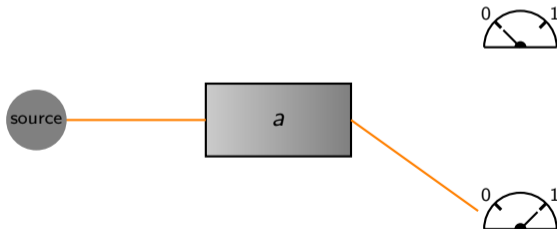
Monogamy relation: $\mathcal{B}(A, B) + \mathcal{B}(A, C) \leq 2R$

Locality of macroscopic averages

Macroscopic measurements

(Microscopic) dichotomic measurement

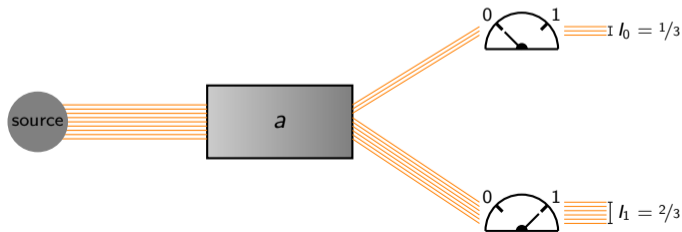
- ▶ A single particle is subject to an interaction a and collides with one of two detectors: outcomes 0 and 1.
- ▶ The interaction is probabilistic: $p(a = x)$, with $x = 0, 1$.



Macroscopic measurements

Macroscopic dichotomic measurement

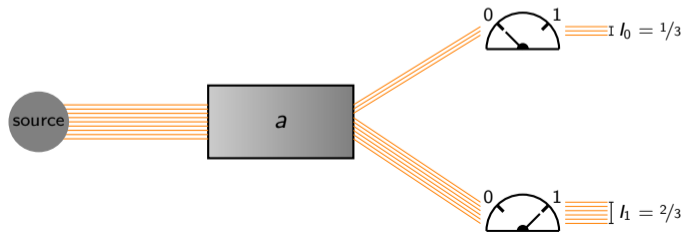
- ▶ Consider beam (or region) of N particles, differently prepared.
- ▶ Subject each particle to the interaction a : the beam gets divided into 2 smaller beams hitting each of the detectors.



Macroscopic measurements

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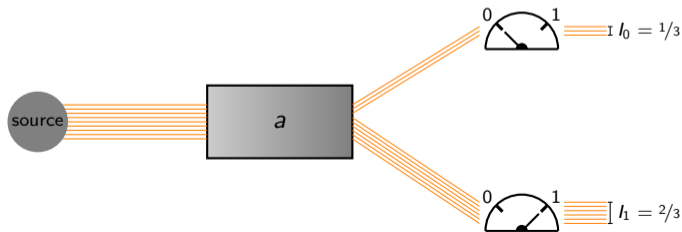
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- ▶ Outcome represented by the intensity of resulting beams:
 $I_1 \in [0, 1]$ proportion of particles hitting the detector 1.



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- ▶ Outcome represented by the intensity of resulting beams:
 $I_1 \in [0, 1]$ proportion of particles hitting the detector 1.
- ▶ We are only concerned with the mean, or **expected value**, of such intensities.



Macroscopic average behaviour

- ▶ This mean intensity can be interpreted as the average behaviour among the N particles:
if we would randomly select one of the particles and subject it to the microscopic measurement a , we would obtain outcome x with probability I_x :

$$I_x = \frac{1}{N} \sum_{i=1}^N p_i(a = x) .$$

- ▶ The situation is analogous to statistical mechanics, where a macrostate arises as an averaging over an extremely large number of microstates, and hence several different microstates can correspond to the same macrostate.

Macroscopic average behaviour: multipartite

▶ **Multipartite macroscopic scenarios**

- ▶ several 'macroscopic' sites consisting of a large number of microscopic sites/particles;
- ▶ several (macro) measurement settings at each site.

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the (mean) values of the macroscopic intensities indicate the behaviour of a randomly chosen tuple of particles: one from each of the beams, or sites.

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- ▶ We'll show that, as long as there are enough particles (microscopic sites) in each macroscopic site, such average macroscopic behaviour is always local no matter which no-signalling model accounts for the underlying microscopic correlations.

Macroscopic average behaviour: (toy) tripartite example

- ▶ We regard sites B and C as forming one 'macroscopic' site, M , and A as forming another.
- ▶ In order to be 'lumped together', B and C must be symmetric/of the same type: the symmetry identifies the measurements $b_0 \sim c_0$ and $b_1 \sim c_1$, giving rise to 'macroscopic' measurements m_0 and m_1 .

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- ▶ Given a (tripartite) empirical model $p(a_i, b_j, c_k = x, y, z)$, the 'macroscopic' average behaviour is a bipartite model (with two macro sites A and M) given by the following average of probabilities of the partial models:

$$p(a_i, m_j = x, y) = \frac{p(a_i, b_j = x, y) + p(a_i, c_j = x, y)}{2}$$

Example: W state

Z and X measurements on the W state:

	000	001	010	011	100	101	110	111
$a_0 b_0 c_0$	9	1	1	1	1	1	1	9
$a_0 b_0 c_1$	8	2	0	2	0	2	8	2
$a_0 b_1 c_0$	8	0	2	2	0	8	2	2
$a_0 b_1 c_1$	4	4	4	0	4	4	4	0
$a_1 b_0 c_0$	8	0	0	8	2	2	2	2
$a_1 b_0 c_1$	4	4	4	4	4	0	4	0
$a_1 b_1 c_0$	4	4	4	4	4	4	0	0
$a_1 b_1 c_1$	0	8	8	0	8	0	0	0

(every entry should be divided by 24)

Example: W state

	00	01	10	11
$a_0 m_0$	10	2	2	10
$a_0 m_1$	8	4	8	4
$a_1 m_0$	8	8	4	4
$a_1 m_1$	8	8	8	0

(every entry should be divided by 24)

This is **local**!

Another example model

	000	001	010	011	100	101	110	111
$a_0 b_0 c_0$	1	1	0	0	0	0	1	1
$a_0 b_0 c_1$	1	1	0	0	0	0	1	1
$a_0 b_1 c_0$	1	1	0	0	0	0	1	1
$a_0 b_1 c_1$	1	1	0	0	0	0	1	1
$a_1 b_0 c_0$	1	1	0	0	0	0	1	1
$a_1 b_0 c_1$	1	1	0	0	0	0	1	1
$a_1 b_1 c_0$	0	0	1	1	1	1	0	0
$a_1 b_1 c_1$	0	0	1	1	1	1	0	0

(every entry should be divided by 4)

Another example model

	00	01	10	11
$a_0 b_0$	2	0	0	2
$a_0 b_1$	2	0	0	2
$a_1 b_0$	2	0	0	2
$a_1 b_1$	0	2	2	0

(divided by 4)

	00	01	10	11
$a_0 c_0$	1	1	1	1
$a_0 c_1$	1	1	1	1
$a_1 c_0$	1	1	1	1
$a_1 c_1$	1	1	1	1

(divided by 4)

maximally non-local

local

	00	01	10	11
$a_0 m_0$	3	1	1	3
$a_0 m_1$	3	1	1	3
$a_1 m_0$	3	1	1	3
$a_1 m_1$	1	3	3	1

(every entry should be divided by 8)

Again, this is **local**!

Connecting monogamy
and macroscopic averages

A simple observation

Consider **any** bipartite Bell inequality $\mathcal{B}(-, -) \leq R$ given by coefficients $\alpha(i, j, x, y)$ and bound R .

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$$\mathcal{B}(A, M) \leq R$$

\Leftrightarrow

$$\sum_{i,j,x,y} \alpha(i, j, x, y) p(a_i, m_j = x, y) \leq R$$

\Leftrightarrow

$$\sum_{i,j,x,y} \alpha(i, j, x, y) \frac{p(a_i, b_j = x, y) + p(a_i, c_j = x, y)}{2} \leq R$$

\Leftrightarrow

$$\sum_{i,j,x,y} \alpha(i, j, x, y) p(a_i, b_j = x, y) + \sum_{i,j,x,y} \alpha(i, j, x, y) p(a_i, c_j = x, y) \leq 2R$$

\Leftrightarrow

$$\mathcal{B}(A, B) + \mathcal{B}(A, C) \leq R$$

A simple observation

Consider **any** bipartite Bell inequality $\mathcal{B}(-, -) \leq R$ given by coefficients $\alpha(i, j, x, y)$ and bound R .

$$\begin{aligned} & \mathcal{B}(A, M) \leq R \\ \Leftrightarrow & \sum_{i,j,x,y} \alpha(i, j, x, y) p(a_i, m_j = x, y) \leq R \\ \Leftrightarrow & \sum_{i,j,x,y} \alpha(i, j, x, y) \frac{p(a_i, b_j = x, y) + p(a_i, c_j = x, y)}{2} \leq R \\ \Leftrightarrow & \sum_{i,j,x,y} \alpha(i, j, x, y) p(a_i, b_j = x, y) + \sum_{i,j,x,y} \alpha(i, j, x, y) p(a_i, c_j = x, y) \leq 2R \\ \Leftrightarrow & \mathcal{B}(A, B) + \mathcal{B}(A, C) \leq R \end{aligned}$$

The average model p_{a_i, m_j} **satisfies the Bell inequality** if and only if in the microscopic model Alice is **monogamous** with respect to violating it with Bob and Charlie.

A simple observation

- ▶ In the two examples above, the average models were local.
- ▶ Equivalently, the examples satisfied the monogamy relation for **any** Bell inequality.

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- ▶ In the two examples above, the average models were local.
- ▶ Equivalently, the examples satisfied the monogamy relation for **any** Bell inequality.
- ▶ This is true for all no-signalling empirical models on the tripartite scenario under consideration, with two measurement settings per site.
- ▶ We now give a structural explanation for this...
- ▶ ... which generalises well beyond this particular scenario.

Vorob'ev's theorem

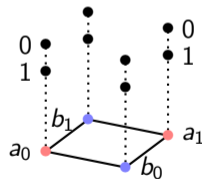
Formalising empirical data

A **measurement scenario** $\mathbf{X} = \langle X, \Sigma, O \rangle$:

- ▶ X – a finite set of measurements
- ▶ $O = (O_x)_{x \in X}$ – for each $x \in X$ a non-empty set of possible outcomes O_x
- ▶ Σ – an abstract simplicial complex on X faces are called the **measurement contexts**

A	B	(0, 0)	(0, 1)	(1, 0)	(1, 1)
a_0	b_0	---	---	---	---
a_0	b_1	---	---	---	---
a_1	b_0	---	---	---	---
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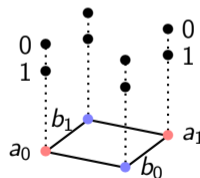
family of finite subsets of X such that:

- ▶ it contains all the singletons:
 $\forall x \in X. \{x\} \in \Sigma.$
- ▶ it is downwards closed: $\sigma \in \Sigma$ and $\tau \subseteq \sigma$ implies $\tau \in \Sigma$.

A	B	(0, 0)	(0, 1)	(1, 0)	(1, 1)
a_0	b_0	---	---	---	---
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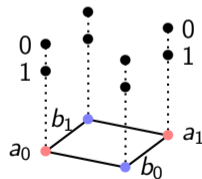
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- ▶ *generalised no-signalling* holds: for any $\sigma, \tau \in \Sigma$, if $\tau \subseteq \sigma$,

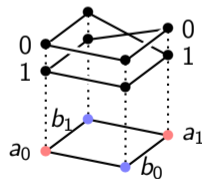
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(i.e. marginals are well-defined)

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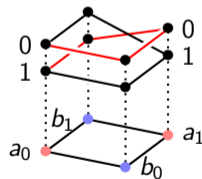
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If no such global distribution exists, the empirical model is **contextual**.

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That is, we can **glue** all the local information together into a global consistent description from which the local information can be recovered.

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The import of Bell's and Kochen–Spekker's theorems is that there are behaviours arising from quantum mechanics that are contextual.

Vorob'ev's theorem

'Consistent families of measures and their extensions'

Vorob'ev, Theory Probab. Appl. 7(2), 1962.

- ▶ In the context of game theory.
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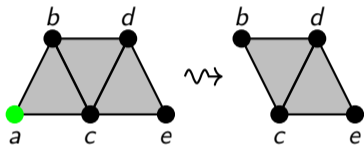
- ▶ Necessary and sufficient condition: **regularity** or **acyclicity**!

Acyclicity

- ▶ Graham reduction step: delete a vertex that belongs to only one maximal face.

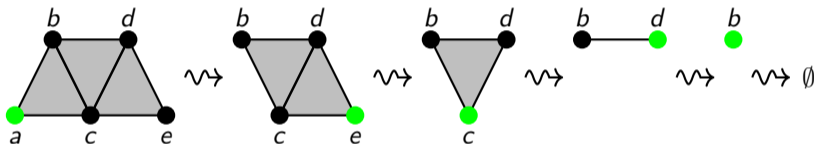
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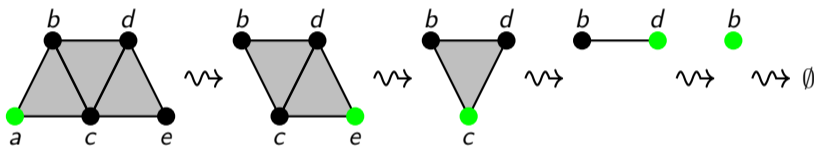
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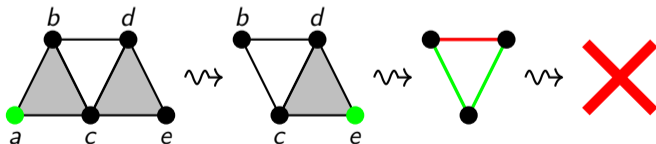


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- ▶ Σ not acyclic: Graham reduction fails.



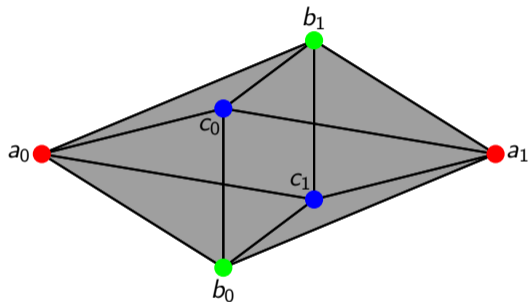
Vorob'ev's theorem

Theorem (Vorob'ev 1962, adapted)

All empirical models on Σ are extendable iff Σ is acyclic

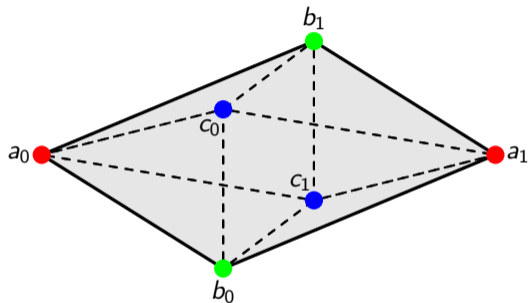
A structural explanation

Structural reason



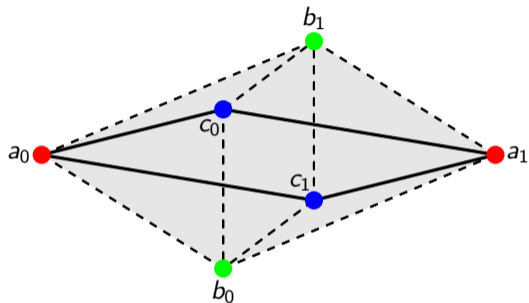
- ▶ Measurement scenario: simplicial complex $\mathfrak{D}_2 * \mathfrak{D}_2 * \mathfrak{D}_2$.

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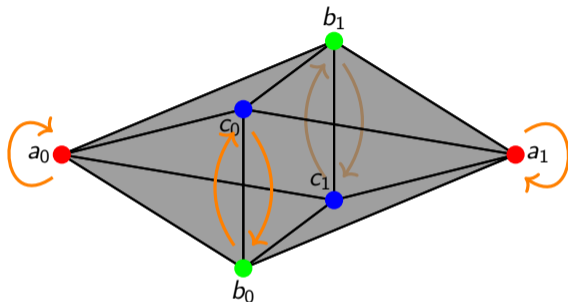
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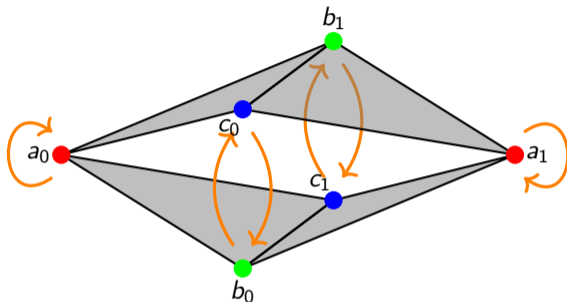
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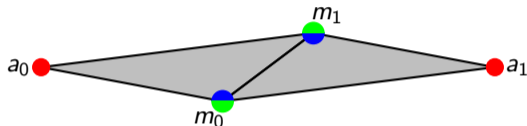
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- ▶ We identify B and C : $b_0 \sim c_0$, $b_1 \sim c_1$.
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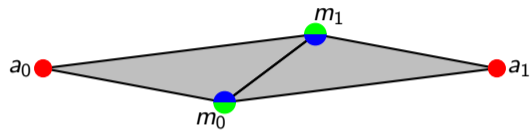
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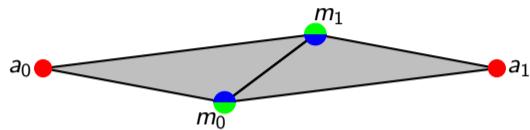


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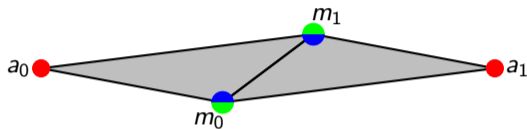


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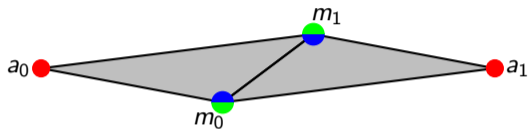
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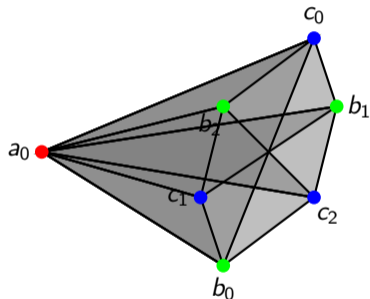
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- ▶ Therefore, no matter from which micro model p_{a_i, b_j, c_k} we start, the averaged macro correlations p_{a_i, m_j} are local.
- ▶ In particular, they satisfy any Bell inequality.
- ▶ Hence, the original tripartite model also satisfies a monogamy relation for any Bell inequality.

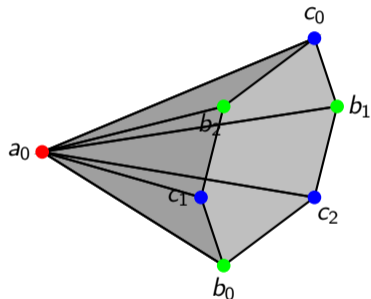
A non-acyclic example

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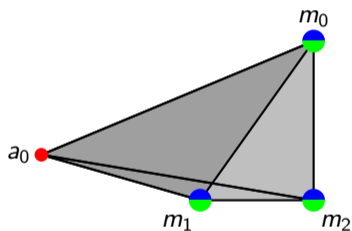
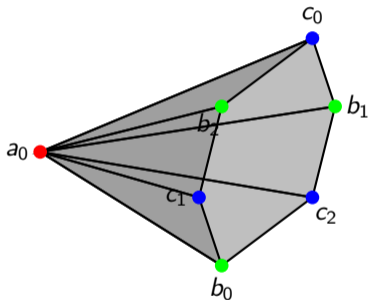
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$$\text{Scenario: } \Sigma_{n, \vec{k}, \vec{r}} := \mathfrak{D}_{k_1}^{(*r_1)} * \dots * \mathfrak{D}_{k_n}^{(*r_n)}$$

- ▶ Symmetry by $S_{r_1} \times \dots \times S_{r_n}$ identifies the copies at each macro site.

$$a_j^{(1)} \sim \dots \sim a_j^{(r_A)} \quad (\forall j \in \{1, \dots, k_A\}),$$

$$b_j^{(1)} \sim \dots \sim a_j^{(r_A)} \quad (\forall j \in \{1, \dots, k_A\}),$$

etc.

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We get monogamy relations

$$\sum_{m_B=1}^{r_B} \sum_{m_C=1}^{r_C} \dots \mathcal{B}(A, B^{(m_B)}, C^{(m_C)}, \dots) \leq r_B r_C \dots R$$

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- ▶ The approach is not restricted to multipartite Bell-type scenarios. More generally, we can apply the same ideas to derive monogamy relations for contextuality inequalities.

Questions...

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