

A First Encounter with General Probabilistic Theories

Rafael Wagner

September 20, 2022

Quantum Theory

Postulates

Structure

General Probabilistic Theories

Classical Probabilistic Theories

Real Quantum Theory

Diagrams

Appendix



Quantum Theory



1. Is quantum theory an island in theory landscape?



1. Is quantum theory an island in theory landscape?
2. Is QT everything?



1. Is quantum theory an island in theory landscape?
2. Is QT everything?
3. Testing nonclassicality experimentally without assuming QT.



1. Is quantum theory an island in theory landscape?
2. Is QT everything?
3. Testing nonclassicality experimentally without assuming QT.
4. It teaches us relevant things about computation.



1. Is quantum theory an island in theory landscape?
2. Is QT everything?
3. Testing nonclassicality experimentally without assuming QT.
4. It teaches us relevant things about computation.
5. Good tools for theoretical investigations. **Diagrams!**



1 To every physical system there corresponds a complex and separable Hilbert space \mathcal{H} .



1 To every physical system there corresponds a complex and separable Hilbert space \mathcal{H} . Every nonzero vector $|\psi\rangle \in \mathcal{H}$ gives a complete description of the state of the system.



- 1 To every physical system there corresponds a complex and separable Hilbert space \mathcal{H} . Every nonzero vector $|\psi\rangle \in \mathcal{H}$ gives a complete description of the state of the system. For each $\lambda \in \mathbb{C}, \lambda \neq 0$, $\lambda|\psi\rangle$ and $|\psi\rangle$ describe the same state.
- 2 The composition of two physical systems described by \mathcal{H}_1 and \mathcal{H}_2



- 1 To every physical system there corresponds a complex and separable Hilbert space \mathcal{H} . Every nonzero vector $|\psi\rangle \in \mathcal{H}$ gives a complete description of the state of the system. For each $\lambda \in \mathbb{C}, \lambda \neq 0$, $\lambda|\psi\rangle$ and $|\psi\rangle$ describe the same state.
- 2 The composition of two physical systems described by \mathcal{H}_1 and \mathcal{H}_2 is a new system described by $\mathcal{H}_1 \otimes \mathcal{H}_2$.
- 3 When no measurement is performed



- 1 To every physical system there corresponds a complex and separable Hilbert space \mathcal{H} . Every nonzero vector $|\psi\rangle \in \mathcal{H}$ gives a complete description of the state of the system. For each $\lambda \in \mathbb{C}, \lambda \neq 0$, $\lambda|\psi\rangle$ and $|\psi\rangle$ describe the same state.
- 2 The composition of two physical systems described by \mathcal{H}_1 and \mathcal{H}_2 is a new system described by $\mathcal{H}_1 \otimes \mathcal{H}_2$.
- 3 When no measurement is performed the change in states ψ is given by $|\psi(t)\rangle = U(t)|\psi\rangle$ for some strongly continuous one-parameter (semi)group $t \mapsto U(t)$.



States as density operators: Given some ensemble of states $\{(|\psi_i\rangle, p_i)\}$,

$$\rho := \sum_i p_i |\psi_i\rangle \langle \psi_i| \quad (1)$$



4 To every measurement outcome E of a measurement $M := \{E\}$ we associated positive operators E such that $\sum_{E \in M} E = \mathbb{1}$.



4 To every measurement outcome E of a measurement $M := \{E\}$ we associated positive operators E such that $\sum_{E \in M} E = \mathbb{1}$.

5 The probability that a result E happens given measurement $M \ni E$ is performed on a state ρ is given by the Born rule $\text{Tr}(\rho E)$.



4 To every measurement outcome E of a measurement $M := \{E\}$ we associated positive operators E such that $\sum_{E \in M} E = \mathbb{1}$.

5 The probability that a result E happens given measurement $M \ni E$ is performed on a state ρ is given by the Born rule $\text{Tr}(\rho E)$.

6 When a measurement $M = \{E\}$ is performed on ρ and outcome E is obtained, there is a discontinuous change towards a new state $\rho' = \frac{\sqrt{E}\rho\sqrt{E}}{\text{Tr}(\rho E)}$.



Foundations: There are many arguments showing that postulates 1,2,3 are more fundamental than postulates 4,5,6.



Foundations: There are many arguments showing that postulates 1,2,3 are more fundamental than postulates 4,5,6.

Physics: Aren't they too mathematical?



Foundations: There are many arguments showing that postulates 1,2,3 are more fundamental than postulates 4,5,6.

Physics: Aren't they too mathematical? Wouldn't be nicer to have something like in special relativity?



Lemma

There is a bijective correspondence between POVM elements E and linear functionals $e_E : \mathcal{B}(\mathcal{H}) \rightarrow \mathbb{C}$ defined by $e_E(\cdot) := \text{Tr}(E(\cdot))$.



Lemma

There is a bijective correspondence between POVM elements E and linear functionals $e_E : \mathcal{B}(\mathcal{H}) \rightarrow \mathbb{C}$ defined by $e_E(\cdot) := \text{Tr}(E(\cdot))$.

Proof.

The inner-product $\langle A, B \rangle_{HS} := \text{Tr}(A^\dagger B)$ makes $\mathcal{B}(\mathcal{H})$ a Hilbert space.



Lemma

There is a bijective correspondence between POVM elements E and linear functionals $e_E : \mathcal{B}(\mathcal{H}) \rightarrow \mathbb{C}$ defined by $e_E(\cdot) := \text{Tr}(E(\cdot))$.

Proof.

The inner-product $\langle A, B \rangle_{HS} := \text{Tr}(A^\dagger B)$ makes $\mathcal{B}(\mathcal{H})$ a Hilbert space. From the Riesz-Fréchet representation theorem



Lemma

There is a bijective correspondence between POVM elements E and linear functionals $e_E : \mathcal{B}(\mathcal{H}) \rightarrow \mathbb{C}$ defined by $e_E(\cdot) := \text{Tr}(E(\cdot))$.

Proof.

The inner-product $\langle A, B \rangle_{HS} := \text{Tr}(A^\dagger B)$ makes $\mathcal{B}(\mathcal{H})$ a Hilbert space. From the Riesz-Fréchet representation theorem to every element $A \in \mathcal{B}(\mathcal{H})$ there is one, and only one, linear functional $a_A \in \mathcal{B}(\mathcal{H})^*$, which is defined by $a_A = \langle A, \cdot \rangle_{HS} = \text{Tr}(A^\dagger(\cdot))$.



Lemma

There is a bijective correspondence between POVM elements E and linear functionals $e_E : \mathcal{B}(\mathcal{H}) \rightarrow \mathbb{C}$ defined by $e_E(\cdot) := \text{Tr}(E(\cdot))$.

Proof.

The inner-product $\langle A, B \rangle_{HS} := \text{Tr}(A^\dagger B)$ makes $\mathcal{B}(\mathcal{H})$ a Hilbert space. From the Riesz-Fréchet representation theorem to every element $A \in \mathcal{B}(\mathcal{H})$ there is one, and only one, linear functional $a_A \in \mathcal{B}(\mathcal{H})^*$, which is defined by $a_A = \langle A, \cdot \rangle_{HS} = \text{Tr}(A^\dagger(\cdot))$. Since E is positive we get the result. □



Lemma

There is a bijective correspondence between POVM elements E and linear functionals $e_E : \mathcal{B}(\mathcal{H}) \rightarrow \mathbb{C}$ defined by $e_E(\cdot) := \text{Tr}(E(\cdot))$.

Proof.

The inner-product $\langle A, B \rangle_{HS} := \text{Tr}(A^\dagger B)$ makes $\mathcal{B}(\mathcal{H})$ a Hilbert space. From the **Riesz-Fréchet representation theorem** to every element $A \in \mathcal{B}(\mathcal{H})$ there is one, and only one, linear functional $a_A \in \mathcal{B}(\mathcal{H})^*$, which is defined by $a_A = \langle A, \cdot \rangle_{HS} = \text{Tr}(A^\dagger(\cdot))$. Since E is positive we get the result. □



$$|\psi\rangle \iff \langle\psi|$$



1. The set of hermitian matrices of dimension d_H over the complex numbers is denoted $\mathbb{H}_{d_H}(\mathbb{C})$ and is isomorphic to $\mathcal{B}(\mathcal{H})^{\mathbb{R}}$ the self-adjoint bounded operators.



1. The set of hermitian matrices of dimension d_H over the complex numbers is denoted $\mathbb{H}_{d_H}(\mathbb{C})$ and is isomorphic to $\mathcal{B}(\mathcal{H})^{\mathbb{R}}$ the self-adjoint bounded operators. States are density matrices,



1. The set of hermitian matrices of dimension d_H over the complex numbers is denoted $\mathbb{H}_{d_H}(\mathbb{C})$ and is isomorphic to $\mathcal{B}(\mathcal{H})^{\mathbb{R}}$ the self-adjoint bounded operators. States are density matrices, meaning that $\rho \in \mathbb{H}_{d_H}(\mathbb{C})$, $\rho \geq 0$, $\text{Tr}(\rho) = 1$.



1. The set of hermitian matrices of dimension d_H over the complex numbers is denoted $\mathbb{H}_{d_H}(\mathbb{C})$ and is isomorphic to $\mathcal{B}(\mathcal{H})^{\mathbb{R}}$ the self-adjoint bounded operators. States are density matrices, meaning that $\rho \in \mathbb{H}_{d_H}(\mathbb{C})$, $\rho \geq 0$, $\text{Tr}(\rho) = 1$. We also have that $\dim(\mathbb{H}_{d_H}(\mathbb{C})) = d_H^2$.
2. $\mathbb{H}_{d_H}(\mathbb{C})^+ := \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : \sigma \geq 0\}$



1. The set of hermitian matrices of dimension d_H over the complex numbers is denoted $\mathbb{H}_{d_H}(\mathbb{C})$ and is isomorphic to $\mathcal{B}(\mathcal{H})^{\mathbb{R}}$ the self-adjoint bounded operators. States are density matrices, meaning that $\rho \in \mathbb{H}_{d_H}(\mathbb{C})$, $\rho \geq 0$, $\text{Tr}(\rho) = 1$. We also have that $\dim(\mathbb{H}_{d_H}(\mathbb{C})) = d_H^2$.
2. $\mathbb{H}_{d_H}(\mathbb{C})^+ := \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : \sigma \geq 0\}$ is a closed



1. The set of hermitian matrices of dimension d_H over the complex numbers is denoted $\mathbb{H}_{d_H}(\mathbb{C})$ and is isomorphic to $\mathcal{B}(\mathcal{H})^{\mathbb{R}}$ the self-adjoint bounded operators. States are density matrices, meaning that $\rho \in \mathbb{H}_{d_H}(\mathbb{C})$, $\rho \geq 0$, $\text{Tr}(\rho) = 1$. We also have that $\dim(\mathbb{H}_{d_H}(\mathbb{C})) = d_H^2$.
2. $\mathbb{H}_{d_H}(\mathbb{C})^+ := \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : \sigma \geq 0\}$ is a closed generating



1. The set of hermitian matrices of dimension d_H over the complex numbers is denoted $\mathbb{H}_{d_H}(\mathbb{C})$ and is isomorphic to $\mathcal{B}(\mathcal{H})^{\mathbb{R}}$ the self-adjoint bounded operators. States are density matrices, meaning that $\rho \in \mathbb{H}_{d_H}(\mathbb{C})$, $\rho \geq 0$, $\text{Tr}(\rho) = 1$. We also have that $\dim(\mathbb{H}_{d_H}(\mathbb{C})) = d_H^2$.
2. $\mathbb{H}_{d_H}(\mathbb{C})^+ := \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : \sigma \geq 0\}$ is a closed generating convex



1. The set of hermitian matrices of dimension d_H over the complex numbers is denoted $\mathbb{H}_{d_H}(\mathbb{C})$ and is isomorphic to $\mathcal{B}(\mathcal{H})^{\mathbb{R}}$ the self-adjoint bounded operators. States are density matrices, meaning that $\rho \in \mathbb{H}_{d_H}(\mathbb{C})$, $\rho \geq 0$, $\text{Tr}(\rho) = 1$. We also have that $\dim(\mathbb{H}_{d_H}(\mathbb{C})) = d_H^2$.
2. $\mathbb{H}_{d_H}(\mathbb{C})^+ := \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : \sigma \geq 0\}$ is a closed generating convex pointed cone.



1. The set of hermitian matrices of dimension d_H over the complex numbers is denoted $\mathbb{H}_{d_H}(\mathbb{C})$ and is isomorphic to $\mathcal{B}(\mathcal{H})^{\mathbb{R}}$ the self-adjoint bounded operators. States are density matrices, meaning that $\rho \in \mathbb{H}_{d_H}(\mathbb{C})$, $\rho \geq 0$, $\text{Tr}(\rho) = 1$. We also have that $\dim(\mathbb{H}_{d_H}(\mathbb{C})) = d_H^2$.
2. $\mathbb{H}_{d_H}(\mathbb{C})^+ := \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : \sigma \geq 0\}$ is a closed generating convex pointed cone.
3. $\Omega := \{\rho \in \mathbb{H}_{d_H}(\mathbb{C})^+ : \text{Tr}(\rho) = 1\}$. $\dim(\Omega) = d_H^2 - 1$.



A matrix σ is positive (semi)definite, and we write $\sigma \geq 0$ iff $\forall v$ we have that $\langle v, \sigma v \rangle \geq 0$.



A matrix σ is positive (semi)definite, and we write $\sigma \geq 0$ iff $\forall v$ we have that $\langle v, \sigma v \rangle \geq 0$.

Lemma

$$\mathbb{H}_{d_H}(\mathbb{C})^+ = \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : sp(\sigma) \subset \mathbb{R}_0^+\}.$$



A matrix σ is positive (semi)definite, and we write $\sigma \geq 0$ iff $\forall v$ we have that $\langle v, \sigma v \rangle \geq 0$.

Lemma

$$\mathbb{H}_{d_H}(\mathbb{C})^+ = \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : sp(\sigma) \subset \mathbb{R}_0^+\}.$$

Proof.

Let $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$ with eigenvectors $\{v_\lambda\}_\lambda$.



A matrix σ is positive (semi)definite, and we write $\sigma \geq 0$ iff $\forall v$ we have that $\langle v, \sigma v \rangle \geq 0$.

Lemma

$$\mathbb{H}_{d_H}(\mathbb{C})^+ = \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : sp(\sigma) \subset \mathbb{R}_0^+\}.$$

Proof.

Let $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$ with eigenvectors $\{v_\lambda\}_\lambda$. Then, if σ is positive, $\langle v_\lambda, \sigma v_\lambda \rangle = \lambda \geq 0, \forall \lambda$.



A matrix σ is positive (semi)definite, and we write $\sigma \geq 0$ iff $\forall v$ we have that $\langle v, \sigma v \rangle \geq 0$.

Lemma

$$\mathbb{H}_{d_H}(\mathbb{C})^+ = \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : \text{sp}(\sigma) \subset \mathbb{R}_0^+\}.$$

Proof.

Let $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$ with eigenvectors $\{v_\lambda\}_\lambda$. Then, if σ is positive, $\langle v_\lambda, \sigma v_\lambda \rangle = \lambda \geq 0, \forall \lambda$. If $\text{sp}(\sigma) \subset \mathbb{R}_0^+$



A matrix σ is positive (semi)definite, and we write $\sigma \geq 0$ iff $\forall v$ we have that $\langle v, \sigma v \rangle \geq 0$.

Lemma

$$\mathbb{H}_{d_H}(\mathbb{C})^+ = \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : \text{sp}(\sigma) \subset \mathbb{R}_0^+\}.$$

Proof.

Let $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$ with eigenvectors $\{v_\lambda\}_\lambda$. Then, if σ is positive, $\langle v_\lambda, \sigma v_\lambda \rangle = \lambda \geq 0, \forall \lambda$. If $\text{sp}(\sigma) \subset \mathbb{R}_0^+$ let $v = \sum_\lambda \alpha_\lambda v_\lambda$



A matrix σ is positive (semi)definite, and we write $\sigma \geq 0$ iff $\forall v$ we have that $\langle v, \sigma v \rangle \geq 0$.

Lemma

$$\mathbb{H}_{d_H}(\mathbb{C})^+ = \{\sigma \in \mathbb{H}_{d_H}(\mathbb{C}) : \text{sp}(\sigma) \subset \mathbb{R}_0^+\}.$$

Proof.

Let $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$ with eigenvectors $\{v_\lambda\}_\lambda$. Then, if σ is positive, $\langle v_\lambda, \sigma v_\lambda \rangle = \lambda \geq 0, \forall \lambda$. If $\text{sp}(\sigma) \subset \mathbb{R}_0^+$ let $v = \sum_\lambda \alpha_\lambda v_\lambda$ then $\langle v, \sigma v \rangle = \sum_{\lambda, \lambda'} \alpha_\lambda \alpha_{\lambda'}^* \lambda \delta_{\lambda, \lambda'} = \sum_\lambda |\alpha_\lambda|^2 \lambda \geq 0$. □





Pointed Convex Cone: (i) (cone)

$$t \geq 0, \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies t\sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+.$$



Pointed Convex Cone: (i) (cone)

$t \geq 0, \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies t\sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+.$ True since
 $\langle v, t\sigma v \rangle = t\langle v, \sigma v \rangle \geq 0$



Pointed Convex Cone: (i) (cone)

$t \geq 0, \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies t\sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+$. True since
 $\langle v, t\sigma v \rangle = t\langle v, \sigma v \rangle \geq 0$ (or also since $\text{sp}(t\sigma) = t\text{sp}(\sigma)$ for any
 $t \in \mathbb{R}_0^+$).



Pointed Convex Cone: (i) (cone)

$t \geq 0, \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies t\sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+$. True since $\langle v, t\sigma v \rangle = t\langle v, \sigma v \rangle \geq 0$ (or also since $\text{sp}(t\sigma) = t\text{sp}(\sigma)$ for any $t \in \mathbb{R}_0^+$).

(ii) (convex) $\sigma, \rho \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies \sigma + \rho \in \mathbb{H}_{d_H}(\mathbb{C})^+$.



Pointed Convex Cone: (i) (cone)

$t \geq 0, \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies t\sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+$. True since $\langle v, t\sigma v \rangle = t\langle v, \sigma v \rangle \geq 0$ (or also since $\text{sp}(t\sigma) = t\text{sp}(\sigma)$ for any $t \in \mathbb{R}_0^+$).

(ii) (convex) $\sigma, \rho \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies \sigma + \rho \in \mathbb{H}_{d_H}(\mathbb{C})^+$. Also true, since $\langle v, (\sigma + \rho)v \rangle = \langle v, \sigma v \rangle + \langle v, \rho v \rangle \geq 0$.



Pointed Convex Cone: (i) (cone)

$t \geq 0, \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies t\sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+$. True since $\langle v, t\sigma v \rangle = t\langle v, \sigma v \rangle \geq 0$ (or also since $\text{sp}(t\sigma) = t\text{sp}(\sigma)$ for any $t \in \mathbb{R}_0^+$).

(ii) (convex) $\sigma, \rho \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies \sigma + \rho \in \mathbb{H}_{d_H}(\mathbb{C})^+$. Also true, since $\langle v, (\sigma + \rho)v \rangle = \langle v, \sigma v \rangle + \langle v, \rho v \rangle \geq 0$.

(iii) (pointed) $\mathbb{H}_{d_H}(\mathbb{C})^+ \cap -\mathbb{H}_{d_H}(\mathbb{C})^+ = \{0\}$.



Pointed Convex Cone: (i) (cone)

$t \geq 0, \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies t\sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+$. True since $\langle v, t\sigma v \rangle = t\langle v, \sigma v \rangle \geq 0$ (or also since $\text{sp}(t\sigma) = t\text{sp}(\sigma)$ for any $t \in \mathbb{R}_0^+$).

(ii) (convex) $\sigma, \rho \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies \sigma + \rho \in \mathbb{H}_{d_H}(\mathbb{C})^+$. Also true, since $\langle v, (\sigma + \rho)v \rangle = \langle v, \sigma v \rangle + \langle v, \rho v \rangle \geq 0$.

(iii) (pointed) $\mathbb{H}_{d_H}(\mathbb{C})^+ \cap -\mathbb{H}_{d_H}(\mathbb{C})^+ = \{0\}$.

$\lambda \in (0, 1), \sigma_1, \sigma_2 \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies \lambda\sigma_1 + (1 - \lambda)\sigma_2 \in \mathbb{H}_{d_H}(\mathbb{C})^+$.



Pointed Convex Cone: (i) (cone)

$t \geq 0, \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies t\sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+$. True since $\langle v, t\sigma v \rangle = t\langle v, \sigma v \rangle \geq 0$ (or also since $\text{sp}(t\sigma) = t\text{sp}(\sigma)$ for any $t \in \mathbb{R}_0^+$).

(ii) (convex) $\sigma, \rho \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies \sigma + \rho \in \mathbb{H}_{d_H}(\mathbb{C})^+$. Also true, since $\langle v, (\sigma + \rho)v \rangle = \langle v, \sigma v \rangle + \langle v, \rho v \rangle \geq 0$.

(iii) (pointed) $\mathbb{H}_{d_H}(\mathbb{C})^+ \cap -\mathbb{H}_{d_H}(\mathbb{C})^+ = \{0\}$.

$\lambda \in (0, 1), \sigma_1, \sigma_2 \in \mathbb{H}_{d_H}(\mathbb{C})^+ \implies \lambda\sigma_1 + (1 - \lambda)\sigma_2 \in \mathbb{H}_{d_H}(\mathbb{C})^+$.

Closed: $\sigma_n \rightarrow \sigma, \langle v, \sigma v \rangle = \langle v, \lim_{n \rightarrow \infty} \sigma_n v \rangle = \lim_{n \rightarrow \infty} \langle v, \sigma_n v \rangle$.



Generating: Every element $\rho \in \mathbb{H}_{d_H}(\mathbb{C})$ is such that $\rho = \rho^+ - \rho^-$ with $\rho^+, \rho^- \in \mathbb{H}_{d_H}(\mathbb{C})^+$.



Real vector-space: Note that any element $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$



Real vector-space: Note that any element $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$ can be described by d_H real numbers in the diagonal



Real vector-space: Note that any element $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$ can be described by d_H real numbers in the diagonal and $\frac{1}{2}d_H(d_H - 1)$ complex numbers above this diagonal



Real vector-space: Note that any element $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$ can be described by d_H real numbers in the diagonal and $\frac{1}{2}d_H(d_H - 1)$ complex numbers above this diagonal (the other half is constrained from transposition).



Real vector-space: Note that any element $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$ can be described by d_H real numbers in the diagonal and $\frac{1}{2}d_H(d_H - 1)$ complex numbers above this diagonal (the other half is constrained from transposition). Therefore we have in total that $2 \cdot \frac{1}{2}d_H(d_H - 1) + d_H = d_H^2 - d_H + d_H = d_H^2$ real numbers for describing these matrices.



Real vector-space: Note that any element $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$ can be described by d_H real numbers in the diagonal and $\frac{1}{2}d_H(d_H - 1)$ complex numbers above this diagonal (the other half is constrained from transposition). Therefore we have in total that $2 \cdot \frac{1}{2}d_H(d_H - 1) + d_H = d_H^2 - d_H + d_H = d_H^2$ real numbers for describing these matrices.

Any matrix $\sigma \in \mathbb{H}_{d_H}(\mathbb{C})$ can then be embedded in vectors in $\mathbb{R}^{d_H^2}$.



Dual cone of a positive cone: For $\mathbb{H}_{d_H}(\mathbb{C})^+$



Dual cone of a positive cone: For $\mathbb{H}_{d_H}(\mathbb{C})^+$ we define the dual cone as $\mathbb{H}_{d_H}(\mathbb{C})^{+*} := \{e \in \mathbb{H}_{d_H}(\mathbb{C})^* : e(\sigma) \geq 0, \forall \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+\}$.



Dual cone of a positive cone: For $\mathbb{H}_{d_H}(\mathbb{C})^+$ we define the dual cone as $\mathbb{H}_{d_H}(\mathbb{C})^{+*} := \{e \in \mathbb{H}_{d_H}(\mathbb{C})^* : e(\sigma) \geq 0, \forall \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+\}$.

Any effect $e = \text{Tr}(E(\cdot))$



Dual cone of a positive cone: For $\mathbb{H}_{d_H}(\mathbb{C})^+$ we define the dual cone as $\mathbb{H}_{d_H}(\mathbb{C})^{+*} := \{e \in \mathbb{H}_{d_H}(\mathbb{C})^* : e(\sigma) \geq 0, \forall \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+\}$.

Any effect $e = \text{Tr}(E(\cdot))$ acting on a positive σ implies $e(\sigma) = \text{Tr}(E\sigma) \geq 0$.



Dual cone of a positive cone: For $\mathbb{H}_{d_H}(\mathbb{C})^+$ we define the dual cone as $\mathbb{H}_{d_H}(\mathbb{C})^{+*} := \{e \in \mathbb{H}_{d_H}(\mathbb{C})^* : e(\sigma) \geq 0, \forall \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+\}$.

Any effect $e = \text{Tr}(E(\cdot))$ acting on a positive σ implies $e(\sigma) = \text{Tr}(E\sigma) \geq 0$. Therefore effects are also elements of a generating convex pointed cone dual to the cone of unnormalized states.



Dual cone of a positive cone: For $\mathbb{H}_{d_H}(\mathbb{C})^+$ we define the dual cone as $\mathbb{H}_{d_H}(\mathbb{C})^{+*} := \{e \in \mathbb{H}_{d_H}(\mathbb{C})^* : e(\sigma) \geq 0, \forall \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+\}$.

Any effect $e = \text{Tr}(E(\cdot))$ acting on a positive σ implies $e(\sigma) = \text{Tr}(E\sigma) \geq 0$. Therefore effects are also elements of a generating convex pointed cone dual to the cone of unnormalized states.

There exists a unique effect $u := \text{Tr}(\mathbb{1}(\cdot))$



Dual cone of a positive cone: For $\mathbb{H}_{d_H}(\mathbb{C})^+$ we define the dual cone as $\mathbb{H}_{d_H}(\mathbb{C})^{+*} := \{e \in \mathbb{H}_{d_H}(\mathbb{C})^* : e(\sigma) \geq 0, \forall \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+\}$.

Any effect $e = \text{Tr}(E(\cdot))$ acting on a positive σ implies $e(\sigma) = \text{Tr}(E\sigma) \geq 0$. Therefore effects are also elements of a generating convex pointed cone dual to the cone of unnormalized states.

There exists a unique effect $u := \text{Tr}(\mathbb{1}(\cdot))$ that says with certainty if any normalized state is present since $u(\sigma) = \text{Tr}(\sigma)$



Dual cone of a positive cone: For $\mathbb{H}_{d_H}(\mathbb{C})^+$ we define the dual cone as $\mathbb{H}_{d_H}(\mathbb{C})^{+*} := \{e \in \mathbb{H}_{d_H}(\mathbb{C})^* : e(\sigma) \geq 0, \forall \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+\}$.

Any effect $e = \text{Tr}(E(\cdot))$ acting on a positive σ implies $e(\sigma) = \text{Tr}(E\sigma) \geq 0$. Therefore effects are also elements of a generating convex pointed cone dual to the cone of unnormalized states.

There exists a unique effect $u := \text{Tr}(\mathbb{1}(\cdot))$ that says with certainty if any normalized state is present since $u(\sigma) = \text{Tr}(\sigma)$ and defines the probability of preparing any state.



Dual cone of a positive cone: For $\mathbb{H}_{d_H}(\mathbb{C})^+$ we define the dual cone as $\mathbb{H}_{d_H}(\mathbb{C})^{+*} := \{e \in \mathbb{H}_{d_H}(\mathbb{C})^* : e(\sigma) \geq 0, \forall \sigma \in \mathbb{H}_{d_H}(\mathbb{C})^+\}$.

Any effect $e = \text{Tr}(E(\cdot))$ acting on a positive σ implies $e(\sigma) = \text{Tr}(E\sigma) \geq 0$. Therefore effects are also elements of a generating convex pointed cone dual to the cone of unnormalized states.

There exists a unique effect $u := \text{Tr}(\mathbb{1}(\cdot))$ that says with certainty if any normalized state is present since $u(\sigma) = \text{Tr}(\sigma)$ and defines the probability of preparing any state. Measurements are sets of effects $M = \{e\}$ such that

$$\sum_{e \in M} e = u \iff \sum_{E \in M} E = \mathbb{1}.$$



Transformations are reversible; for any unitary operator describing the evolution U there exists the inverse operator $U^{-1} = U^\dagger$.



General Probabilistic Theories



Definition

A GPT system is a tuple $(\mathcal{V}, \mathcal{V}_+, \mathcal{E}_{\mathcal{V}}, \mathcal{I}, u)$



Definition

A GPT system is a tuple $(\mathcal{V}, \mathcal{V}_+, \mathcal{E}_{\mathcal{V}}, \mathcal{I}, u)$ where \mathcal{V} is a finite-dimensional real vector space,



Definition

A GPT system is a tuple $(\mathcal{V}, \mathcal{V}_+, \mathcal{E}_{\mathcal{V}}, \mathcal{I}, u)$ where \mathcal{V} is a finite-dimensional real vector space, \mathcal{V}_+ is a closed convex pointed cone generating \mathcal{V} .



Definition

A GPT system is a tuple $(\mathcal{V}, \mathcal{V}_+, \mathcal{E}_{\mathcal{V}}, \mathcal{I}, u)$ where \mathcal{V} is a finite-dimensional real vector space, \mathcal{V}_+ is a closed convex pointed cone generating \mathcal{V} . $\mathcal{E}_{\mathcal{V}}$ is a generating subset of the dual cone $\mathcal{V}_+^* := \{e \in \mathcal{V}^* : e(\omega) \geq 0, \forall \omega \in \mathcal{V}\}$.



Definition

A GPT system is a tuple $(\mathcal{V}, \mathcal{V}_+, \mathcal{E}_{\mathcal{V}}, \mathcal{I}, u)$ where \mathcal{V} is a finite-dimensional real vector space, \mathcal{V}_+ is a closed convex pointed cone generating \mathcal{V} . $\mathcal{E}_{\mathcal{V}}$ is a generating subset of the dual cone $\mathcal{V}_+^* := \{e \in \mathcal{V}^* : e(\omega) \geq 0, \forall \omega \in \mathcal{V}\}$. u is a unique element in $\mathcal{E}_{\mathcal{V}}$, called the order-unit, or also the causal effect.



Definition

A GPT system is a tuple $(\mathcal{V}, \mathcal{V}_+, \mathcal{E}_{\mathcal{V}}, \mathcal{T}, u)$ where \mathcal{V} is a finite-dimensional real vector space, \mathcal{V}_+ is a closed convex pointed cone generating \mathcal{V} . $\mathcal{E}_{\mathcal{V}}$ is a generating subset of the dual cone $\mathcal{V}_+^* := \{e \in \mathcal{V}^* : e(\omega) \geq 0, \forall \omega \in \mathcal{V}\}$. u is a unique element in $\mathcal{E}_{\mathcal{V}}$, called the order-unit, or also the causal effect. Finally, the set \mathcal{T} represents all the linear maps $\mathcal{T} \ni T : \mathcal{V} \rightarrow \mathcal{V}$ such that $T(\mathcal{V}_+) \subseteq \mathcal{V}_+$.



Definition

A composition between two GPT systems

$$A = (A, A_+, E_A, \mathcal{T}_A, u_A) \text{ and } B = (B, B_+, E_B, \mathcal{T}_B, u_B)$$



Definition

A composition between two GPT systems

$A = (A, A_+, E_A, \mathcal{T}_A, u_A)$ and $B = (B, B_+, E_B, \mathcal{T}_B, u_B)$ is a new GPT system AB



Definition

A composition between two GPT systems

$A = (A, A_+, E_A, \mathcal{T}_A, u_A)$ and $B = (B, B_+, E_B, \mathcal{T}_B, u_B)$ is a new GPT system AB together with two bilinear maps



Definition

A composition between two GPT systems

$A = (A, A_+, E_A, \mathcal{T}_A, u_A)$ and $B = (B, B_+, E_B, \mathcal{T}_B, u_B)$ is a new GPT system AB together with two bilinear maps $A \times B \rightarrow AB$ and $A^* \times B^* \rightarrow (AB)^*$ both denoted by \circ , such that

1. Products of normalized states are again normalized states:
 $\omega_A \in \Omega_A, \omega_B \in \Omega_B \implies \omega_{AB} \in \Omega_{AB}.$



Definition

A composition between two GPT systems

$A = (A, A_+, E_A, \mathcal{T}_A, u_A)$ and $B = (B, B_+, E_B, \mathcal{T}_B, u_B)$ is a new GPT system AB together with two bilinear maps $A \times B \rightarrow AB$ and $A^* \times B^* \rightarrow (AB)^*$ both denoted by \circ , such that

1. Products of normalized states are again normalized states:
 $\omega_A \in \Omega_A, \omega_B \in \Omega_B \implies \omega_{AB} \in \Omega_{AB}.$
2. Same for valid effects in E_A and E_B with $u_A \circ u_B \leq u_{AB}$ in particular (local measurements don't lead to probability larger than 1).



Definition

A composition between two GPT systems

$A = (A, A_+, E_A, \mathcal{T}_A, u_A)$ and $B = (B, B_+, E_B, \mathcal{T}_B, u_B)$ is a new GPT system AB together with two bilinear maps $A \times B \rightarrow AB$ and $A^* \times B^* \rightarrow (AB)^*$ both denoted by \circ , such that

1. Products of normalized states are again normalized states:
 $\omega_A \in \Omega_A, \omega_B \in \Omega_B \implies \omega_{AB} \in \Omega_{AB}.$
2. Same for valid effects in E_A and E_B with $u_A \circ u_B \leq u_{AB}$ in particular (local measurements don't lead to probability larger than 1).
3. Local measurements on product states yield statistically independent results. $e_A \circ e_B(\omega_A \circ \omega_B) = e_A(\omega_A) \cdot e_B(\omega_B).$



Definition

A composition between two GPT systems

$A = (A, A_+, E_A, \mathcal{T}_A, u_A)$ and $B = (B, B_+, E_B, \mathcal{T}_B, u_B)$ is a new GPT system AB together with two bilinear maps $A \times B \rightarrow AB$ and $A^* \times B^* \rightarrow (AB)^*$ both denoted by \circ , such that

1. Products of normalized states are again normalized states:
 $\omega_A \in \Omega_A, \omega_B \in \Omega_B \implies \omega_{AB} \in \Omega_{AB}.$
2. Same for valid effects in E_A and E_B with $u_A \circ u_B \leq u_{AB}$ in particular (local measurements don't lead to probability larger than 1).
3. Local measurements on product states yield statistically independent results. $e_A \circ e_B(\omega_A \circ \omega_B) = e_A(\omega_A) \cdot e_B(\omega_B).$
4. $e_{AB}(\omega_A \circ (\cdot)), e_{AB}((\cdot) \circ \phi_B) \in E_{AB}, \forall \phi_B, \omega_A$ normalized
 Similar construction for states.



We have seen that Quantum Theory is a GPT.



We have seen that Quantum Theory is a GPT. It is actually a GPT that has much more structure.



We have seen that Quantum Theory is a GPT. It is actually a GPT that has much more structure. For instance, it satisfies *more properties* that are *not* universal for GPTs.



We have seen that Quantum Theory is a GPT. It is actually a GPT that has much more structure. For instance, it satisfies *more properties* that are *not* universal for GPTs.

No-restriction:



We have seen that Quantum Theory is a GPT. It is actually a GPT that has much more structure. For instance, it satisfies *more properties* that are *not* universal for GPTs.

No-restriction: The set of physical effects $\mathcal{E}_{\mathcal{V}} = \mathcal{V}_+^*$, i.e., every mathematical effect is also a physically implementable one.



We have seen that Quantum Theory is a GPT. It is actually a GPT that has much more structure. For instance, it satisfies *more properties* that are *not* universal for GPTs.

No-restriction: The set of physical effects $\mathcal{E}_{\mathcal{V}} = \mathcal{V}_+^*$, i.e., every mathematical effect is also a physically implementable one.

Strong self-duality:



We have seen that Quantum Theory is a GPT. It is actually a GPT that has much more structure. For instance, it satisfies *more properties* that are *not* universal for GPTs.

No-restriction: The set of physical effects $\mathcal{E}_{\mathcal{V}} = \mathcal{V}_+^*$, i.e., every mathematical effect is also a physically implementable one.

Strong self-duality: $\mathbb{H}_{d_H}(\mathbb{C})^{+*} = \mathbb{H}_{d_H}(\mathbb{C})^+$.



We have seen that Quantum Theory is a GPT. It is actually a GPT that has much more structure. For instance, it satisfies *more properties* that are *not* universal for GPTs.

No-restriction: The set of physical effects $\mathcal{E}_{\mathcal{V}} = \mathcal{V}_+^*$, i.e., every mathematical effect is also a physically implementable one.

Strong self-duality: $\mathbb{H}_{d_H}(\mathbb{C})^{+*} = \mathbb{H}_{d_H}(\mathbb{C})^+$.

Local Tomography:



We have seen that Quantum Theory is a GPT. It is actually a GPT that has much more structure. For instance, it satisfies *more properties* that are *not* universal for GPTs.

No-restriction: The set of physical effects $\mathcal{E}_{\mathcal{V}} = \mathcal{V}_+^*$, i.e., every mathematical effect is also a physically implementable one.

Strong self-duality: $\mathbb{H}_{d_H}(\mathbb{C})^{+*} = \mathbb{H}_{d_H}(\mathbb{C})^+$.

Local Tomography: Any state of a composite system is completely specified by local operations.



We have seen that Quantum Theory is a GPT. It is actually a GPT that has much more structure. For instance, it satisfies *more properties* that are *not* universal for GPTs.

No-restriction: The set of physical effects $\mathcal{E}_{\mathcal{V}} = \mathcal{V}_+^*$, i.e., every mathematical effect is also a physically implementable one.

Strong self-duality: $\mathbb{H}_{d_H}(\mathbb{C})^{+*} = \mathbb{H}_{d_H}(\mathbb{C})^+$.

Local Tomography: Any state of a composite system is completely specified by local operations.





Definition

A CPT system represents a classical random variable taking n different values and is described by the following tuple $(\mathbb{R}^n, \mathbb{R}_+^n, \mathbb{R}_+^n, S_n, \vec{1})$.



Definition

A CPT system represents a classical random variable taking n different values and is described by the following tuple

$$(\mathbb{R}^n, \mathbb{R}_+^n, \mathbb{R}_+^n, S_n, \vec{1}).$$

The identity selects the normalized classical states

$$\vec{1} \cdot \vec{p} = \sum_{i=1}^n p_i = 1.$$



Definition

A CPT system represents a classical random variable taking n different values and is described by the following tuple

$$(\mathbb{R}^n, \mathbb{R}_+^n, \mathbb{R}_+^n, S_n, \vec{1}).$$

The identity selects the normalized classical states

$\vec{1} \cdot \vec{p} = \sum_{i=1}^n p_i = 1$. The set of all normalized states forms a simplex.



Definition

A CPT system represents a classical random variable taking n different values and is described by the following tuple

$$(\mathbb{R}^n, \mathbb{R}_+^n, \mathbb{R}_+^n, S_n, \vec{1}).$$

The identity selects the normalized classical states

$\vec{1} \cdot \vec{p} = \sum_{i=1}^n p_i = 1$. The set of all normalized states forms a simplex.

$$\Delta^{(n)} := \{\vec{p} \in \mathbb{R}_+^n : \sum_i p_i = 1\} \quad (2)$$



Definition

A CPT system represents a classical random variable taking n different values and is described by the following tuple

$$(\mathbb{R}^n, \mathbb{R}_+^n, \mathbb{R}_+^n, S_n, \vec{1}).$$

The identity selects the normalized classical states

$\vec{1} \cdot \vec{p} = \sum_{i=1}^n p_i = 1$. The set of all normalized states forms a simplex.

$$\Delta^{(n)} := \{\vec{p} \in \mathbb{R}_+^n : \sum_i p_i = 1\} \quad (2)$$

Any effect \vec{e} acting on a state \vec{p} is just the inner product

$$\vec{e} \cdot \vec{p} = \sum_{i=1}^n e_i p_i.$$



Definition

A CPT system represents a classical random variable taking n different values and is described by the following tuple

$$(\mathbb{R}^n, \mathbb{R}_+^n, \mathbb{R}_+^n, S_n, \vec{1}).$$

The identity selects the normalized classical states

$\vec{1} \cdot \vec{p} = \sum_{i=1}^n p_i = 1$. The set of all normalized states forms a simplex.

$$\Delta^{(n)} := \{\vec{p} \in \mathbb{R}_+^n : \sum_i p_i = 1\} \quad (2)$$

Any effect \vec{e} acting on a state \vec{p} is just the inner product

$\vec{e} \cdot \vec{p} = \sum_{i=1}^n e_i p_i$. If \vec{e} is also a state, $\vec{e} = \vec{p}_e$ then

$$\vec{p}_e \cdot \vec{p} = \sum_{i=1}^n p_{e_i} p_i.$$



Definition

A CPT system represents a classical random variable taking n different values and is described by the following tuple

$$(\mathbb{R}^n, \mathbb{R}_+^n, \mathbb{R}_+^n, S_n, \vec{1}).$$

The identity selects the normalized classical states

$\vec{1} \cdot \vec{p} = \sum_{i=1}^n p_i = 1$. The set of all normalized states forms a simplex.

$$\Delta^{(n)} := \{\vec{p} \in \mathbb{R}_+^n : \sum_i p_i = 1\} \quad (2)$$

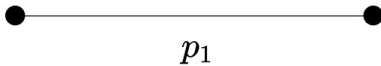
Any effect \vec{e} acting on a state \vec{p} is just the inner product

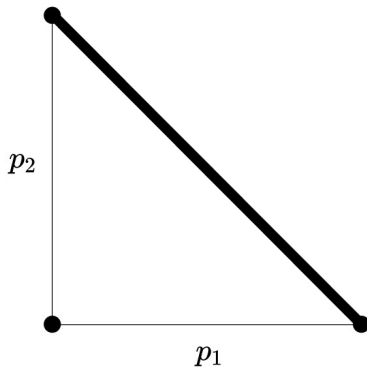
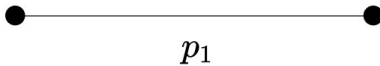
$\vec{e} \cdot \vec{p} = \sum_{i=1}^n e_i p_i$. If \vec{e} is also a state, $\vec{e} = \vec{p}_e$ then

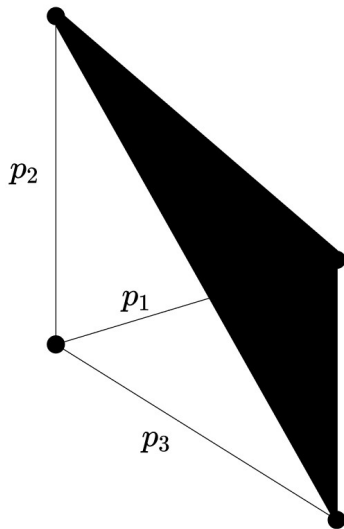
$\vec{p}_e \cdot \vec{p} = \sum_{i=1}^n p_{e_i} p_i$. The normalized effects are $\Delta^{(n)*}$.











What if the following happens...

26



Suppose we have any GPT system A .



Suppose we have any GPT system A . Suppose that we can find a linear map ι such that $\Omega_A \ni \omega \xrightarrow{\iota} \iota(\omega) \in \Delta^{(n)}$.



Suppose we have any GPT system A . Suppose that we can find a linear map ι such that $\Omega_A \ni \omega \mapsto \iota(\omega) \in \Delta^{(n)}$. Suppose also that there is another map κ such that $E_A \ni e \mapsto \kappa(e) \in \Delta^{(n)*}$ that satisfy the following:



Suppose we have any GPT system A . Suppose that we can find a linear map ι such that $\Omega_A \ni \omega \mapsto \iota(\omega) \in \Delta^{(n)}$. Suppose also that there is another map κ such that $E_A \ni e \mapsto \kappa(e) \in \Delta^{(n)*}$ that satisfy the following:

$$\kappa(e) \cdot \iota(\omega) = e(\omega) \tag{3}$$

$$\kappa(u_A) = \vec{1}_n \tag{4}$$



Suppose we have any GPT system A . Suppose that we can find a linear map ι such that $\Omega_A \ni \omega \mapsto \iota(\omega) \in \Delta^{(n)}$. Suppose also that there is another map κ such that $E_A \ni e \mapsto \kappa(e) \in \Delta^{(n)*}$ that satisfy the following:

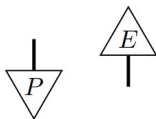
$$\kappa(e) \cdot \iota(\omega) = e(\omega) \tag{3}$$

$$\kappa(u_A) = \vec{1}_n \tag{4}$$

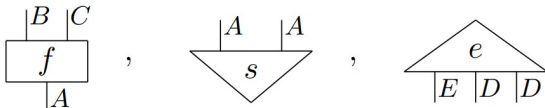
That is noncontextuality in GPT sense!



Nice feature of GPTs: Diagrammatic language!

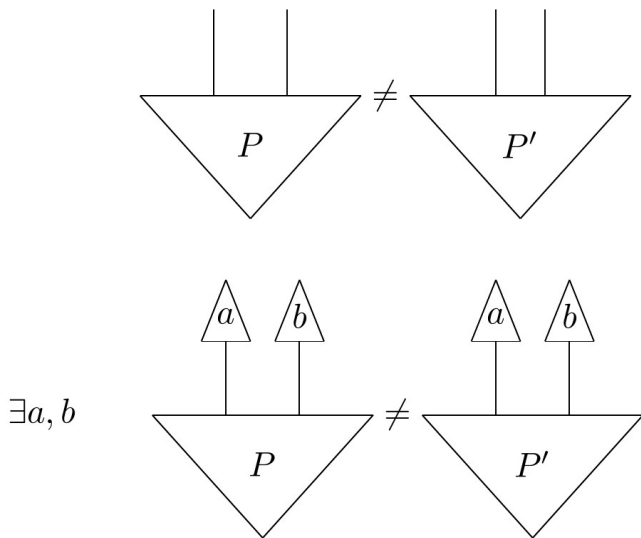


$$f \otimes g \quad := \quad \begin{array}{c} | \\ \boxed{f} \\ | \end{array} \quad \begin{array}{c} | \\ \boxed{g} \\ | \end{array}$$



$$\forall E, E', P, P' \quad \begin{array}{c} \triangle E \quad \triangle E' \\ | \quad | \\ \boxed{T} \\ | \quad | \\ \nabla P \quad \nabla P' \end{array} = \begin{array}{c} \triangle E \quad \triangle E' \\ | \quad | \\ \boxed{T'} \\ | \quad | \\ \nabla P \quad \nabla P' \end{array}$$
$$\Updownarrow$$
$$\begin{array}{c} | \quad | \\ \boxed{T} \\ | \quad | \end{array} = \begin{array}{c} | \quad | \\ \boxed{T'} \\ | \quad | \end{array}.$$





Less Algebra and More Geometry:



Less Algebra and More Geometry: Let's discuss $\mathbb{R}QT$ as a physical theory. We start with the real quantum bit (rebit).



Less Algebra and More Geometry: Let's discuss $\mathbb{R}QT$ as a physical theory. We start with the real quantum bit (rebit).

$$\sigma \in \mathbb{H}_d(\mathbb{R}), \dim(\mathbb{H}_d(\mathbb{R})) = d + \frac{1}{2}d(d-1) = \frac{d^2-d+2d}{2} = \frac{d(d+1)}{2}.$$



Less Algebra and More Geometry: Let's discuss $\mathbb{R}QT$ as a physical theory. We start with the real quantum bit (rebit).

$$\sigma \in \mathbb{H}_d(\mathbb{R}), \dim(\mathbb{H}_d(\mathbb{R})) = d + \frac{1}{2}d(d-1) = \frac{d^2-d+2d}{2} = \frac{d(d+1)}{2}.$$

Therefore the real bits $d = 2$ are embeddable in $2(2+1)/2 = 3$ real vector space.



Less Algebra and More Geometry: Let's discuss $\mathbb{R}QT$ as a physical theory. We start with the real quantum bit (rebit).

$$\sigma \in \mathbb{H}_d(\mathbb{R}), \dim(\mathbb{H}_d(\mathbb{R})) = d + \frac{1}{2}d(d-1) = \frac{d^2-d+2d}{2} = \frac{d(d+1)}{2}.$$

Therefore the real bits $d = 2$ are embeddable in $2(2+1)/2 = 3$ real vector space.

Here is a (choice) basis of the space $\frac{1}{2}\{\mathbb{1}, X, Z\}$.

$$\sigma = a_0 \frac{\mathbb{1}}{2} + a_1 \frac{1}{2}X + a_2 \frac{1}{2}Z \sim \begin{pmatrix} a_0 \\ a_1 \\ a_2 \end{pmatrix}$$



Normalized states:



Normalized states: $1 = u(\sigma) = \text{Tr}(\mathbb{1}\sigma) =$
 $a_0 \text{Tr}(\mathbb{1})/2 + a_1 \text{Tr}(X)/2 + a_2 \text{Tr}(Z)/2 = a_0.$



Normalized states: $1 = u(\sigma) = \text{Tr}(\mathbb{1}\sigma) =$
 $a_0 \text{Tr}(\mathbb{1})/2 + a_1 \text{Tr}(X)/2 + a_2 \text{Tr}(Z)/2 = a_0.$

Positivity:



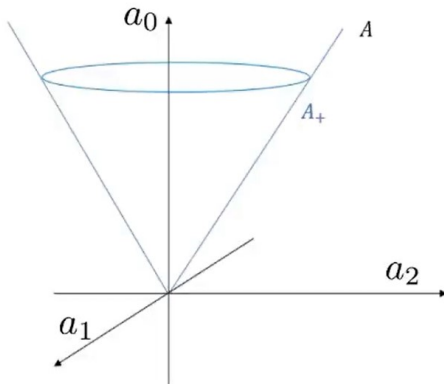
Normalized states: $1 = u(\sigma) = \text{Tr}(\mathbb{1}\sigma) =$
 $a_0 \text{Tr}(\mathbb{1})/2 + a_1 \text{Tr}(X)/2 + a_2 \text{Tr}(Z)/2 = a_0.$

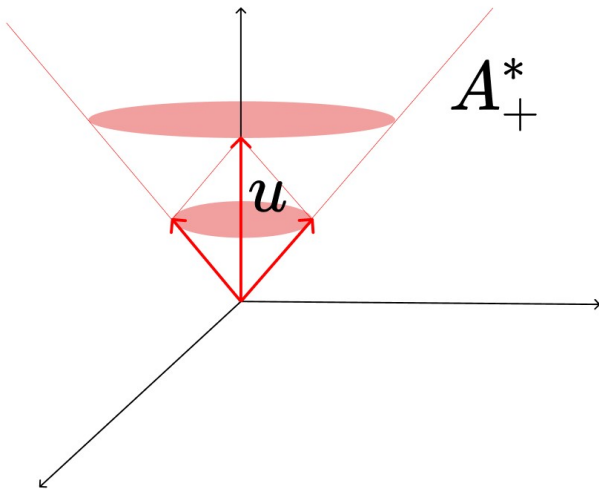
Positivity: $a_1^2 + a_2^2 \leq a_0^2.$



Normalized states: $1 = u(\sigma) = \text{Tr}(1\sigma) =$
 $a_0 \text{Tr}(1)/2 + a_1 \text{Tr}(X)/2 + a_2 \text{Tr}(Z)/2 = a_0.$

Positivity: $a_1^2 + a_2^2 \leq a_0^2.$





Dim. of rebits: $\dim(\mathbb{H}_2) = \frac{d(d+1)}{2} \stackrel{d=2}{=} 3.$



Dim. of rebits: $\dim(\mathbb{H}_2) = \frac{d(d+1)}{2} \stackrel{d=2}{=} 3.$

Dim. of two rebits: $\dim(\mathbb{H}_4) = \frac{2d(2d+1)}{2} \stackrel{d=2}{=} 10.$



Dim. of rebits: $\dim(\mathbb{H}_2) = \frac{d(d+1)}{2} \stackrel{d=2}{=} 3.$

Dim. of two rebits: $\dim(\mathbb{H}_4) = \frac{2d(2d+1)}{2} \stackrel{d=2}{=} 10.$

Dim. of two q-rebits: $\dim(\mathbb{H}_2 \otimes \mathbb{H}_2) = \dim(\mathbb{H}_2) \dim(\mathbb{H}_2) = 9.$



A GPT composition AB is local tomographic if, and only if, the vector spaces compose as the tensor product $AB = A \otimes B$.

Proof.

$$AB = A \otimes B \implies d_{AB} = d_A d_B.$$



A GPT composition AB is local tomographic if, and only if, the vector spaces compose as the tensor product $AB = A \otimes B$.

Proof.

$AB = A \otimes B \implies d_{AB} = d_A d_B$. This implies that the set of all product effects span the vector space $\text{span}(E_{AB})$.



A GPT composition AB is local tomographic if, and only if, the vector spaces compose as the tensor product $AB = A \otimes B$.

Proof.

$AB = A \otimes B \implies d_{AB} = d_A d_B$. This implies that the set of all product effects span the vector space $\text{span}(E_{AB})$. (Any subset A of the dual B^* is separating for B iff A span B^* .)



A GPT composition AB is local tomographic if, and only if, the vector spaces compose as the tensor product $AB = A \otimes B$.

Proof.

$AB = A \otimes B \implies d_{AB} = d_A d_B$. This implies that the set of all product effects span the vector space $\text{span}(E_{AB})$. (Any subset A of the dual B^* is separating for B iff A span B^* .) We have then that the set of product effects is separating, and hence local tomography is satisfied. \square



Real QT is *different* than QT.



Real QT is *different* than QT.

$$Y \otimes Y = \begin{pmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix}$$



Real QT is *different* than QT.

$$Y \otimes Y = \begin{pmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix}$$

$$\frac{\mathbb{1} \otimes \mathbb{1} + Y \otimes Y}{4} \stackrel{QT}{=} \frac{1}{2} (|+i\rangle\langle+i| \otimes |+i\rangle\langle+i| + |-i\rangle\langle-i| \otimes |-i\rangle\langle-i|)$$



Real QT is *different* than QT.

$$Y \otimes Y = \begin{pmatrix} 0 & 0 & 0 & -1 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ -1 & 0 & 0 & 0 \end{pmatrix}$$

$$\frac{\mathbb{1} \otimes \mathbb{1} + Y \otimes Y}{4} \stackrel{QT}{=} \frac{1}{2} (|+i\rangle\langle+i| \otimes |+i\rangle\langle+i| + |-i\rangle\langle-i| \otimes |-i\rangle\langle-i|)$$

Some separable states in *QT* become entangled in *ℝQT*!



Wed 12th May 2021 at 14.00

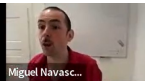
Angelos Bampounis (INL / U Minho)

Journal club: Quantum physics needs complex numbers

Abstract

Slides





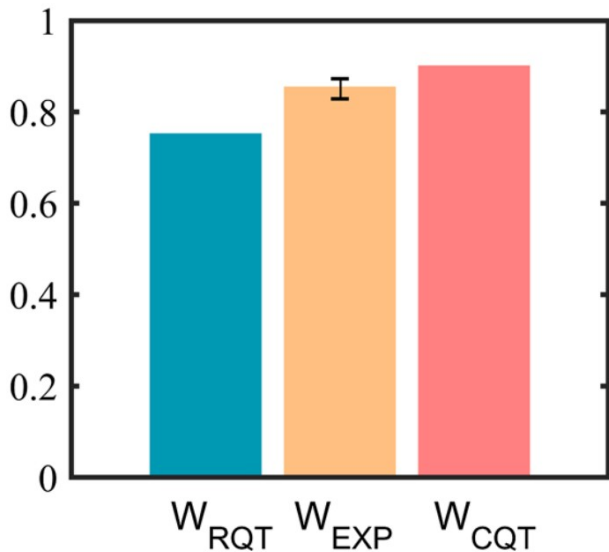
The world is not real

M. O. Renau, D. Trillo, M. Weilenmann, A. Tavakoli,
N. Gisin, A. Acín and M. Navascués

Institute for Quantum Optics and Quantum Information (IQOQI), Vienna

Nature 600, 625–629 (2021).





1. **Local Tomography:** Local operations suffice to learn about states.
2. **No-restriction:** Every mathematically possible effect is also physically possible. $E_A = A_+^*$
3. **Strong self-duality:** There exists an inner-product making effect and state space identical.



Wed 9th December 2020 at 14.00

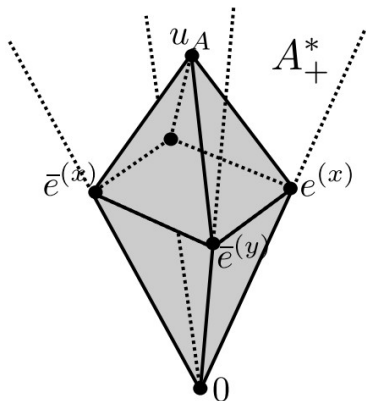
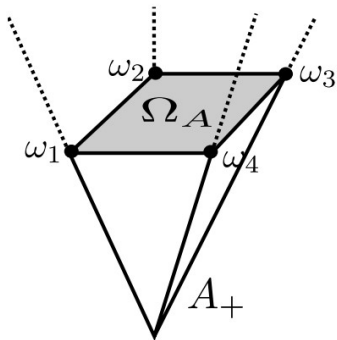
Michael Oliveira (U Minho)

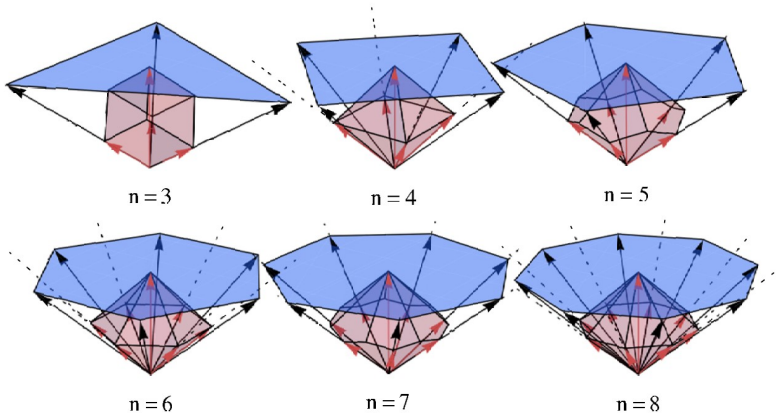
From EPR to Bell, with a description on the correlation polytopes

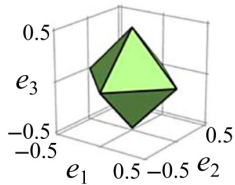
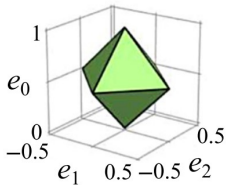
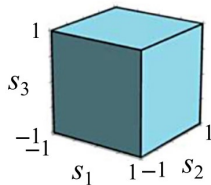
[Abstract](#)

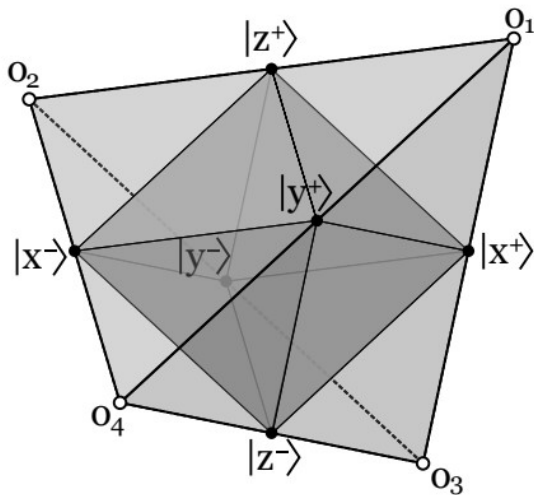
[Slides](#)

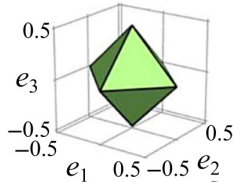
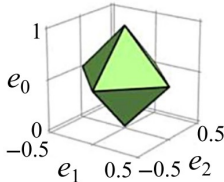
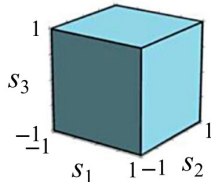
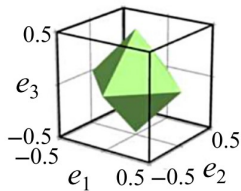
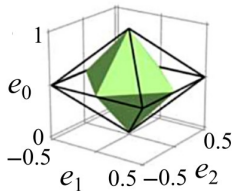
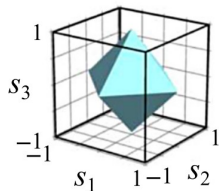


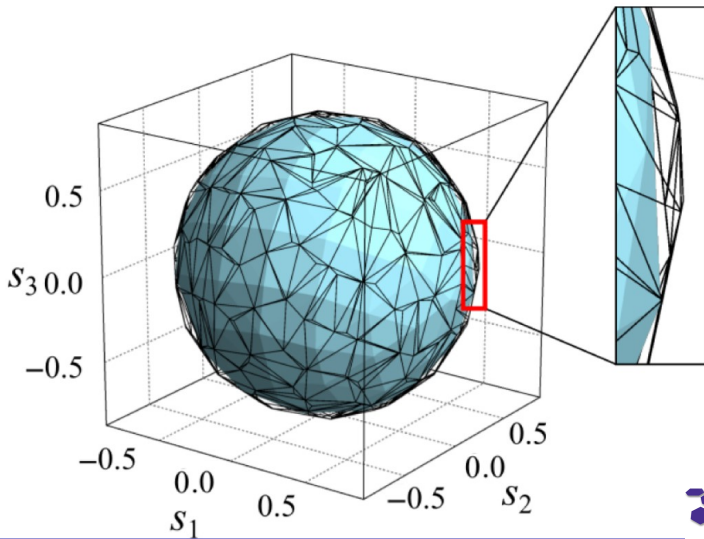


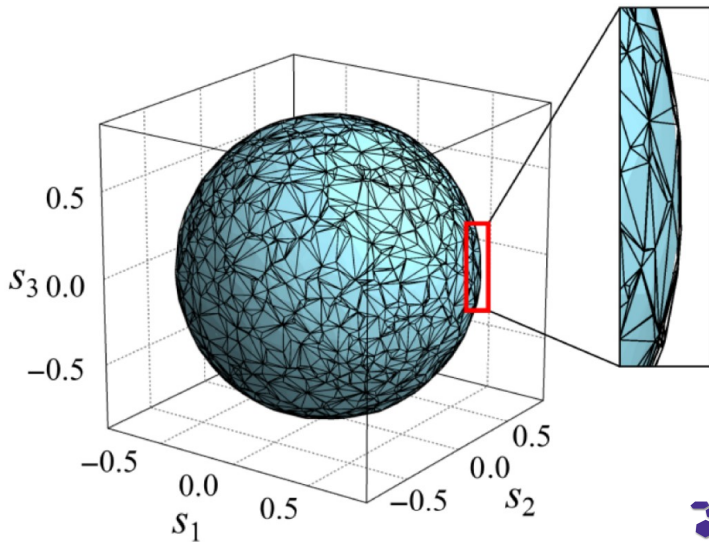












Causality Principle:



Causality Principle: The probability of preparing states is independent of future measurements.



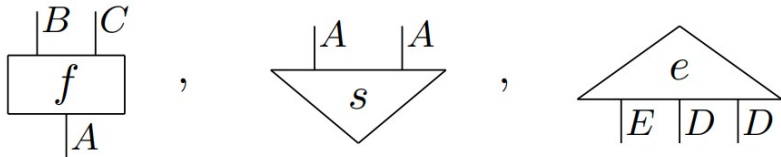
Causality Principle: The probability of preparing states is independent of future measurements.

Theorem (Pavia group)

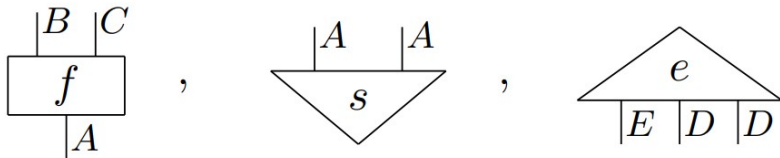
A GPT is causal if, and only if there exists one and only one deterministic effect u .



Causal Tomographically Local GPTs have a diagrammatic language.



Causal Tomographically Local GPTs have a diagrammatic language.



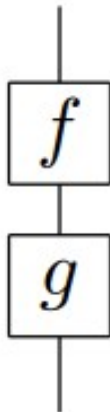
Mathematically, the structure is of a **strict symmetric monoidal category**.



$$\left| A \right\rangle , \quad \left| B \right\rangle \left| C \right\rangle , \quad \left| D \right\rangle \left| D \right\rangle \left| E \right\rangle , \quad \dots$$

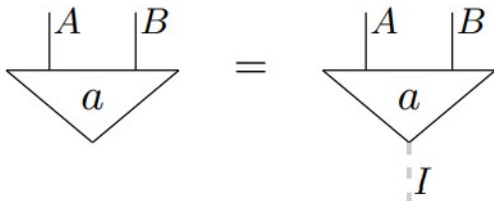


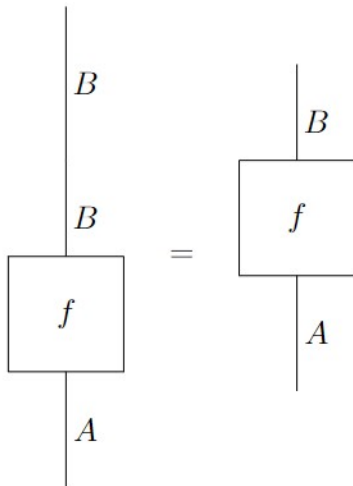
$$f \circ g \quad :=$$



$$\mathbb{1}_A := \begin{array}{c} | A \\ \hline | A \end{array}$$







$$f \otimes g \quad := \quad \begin{array}{c} | \\ \boxed{f} \\ | \end{array} \quad \begin{array}{c} | \\ \boxed{g} \\ | \end{array}$$



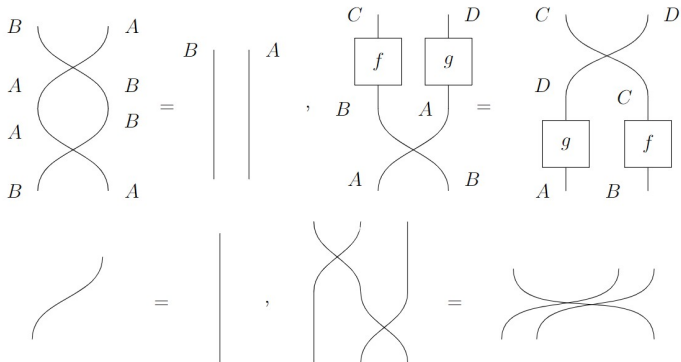
$$(f \otimes_1 g) \circ (h \otimes_1 w) = \begin{array}{c} \text{---} \\ \boxed{f} \\ \text{---} \\ \boxed{h} \\ \text{---} \end{array} \quad \begin{array}{c} \text{---} \\ \boxed{g} \\ \text{---} \\ \boxed{w} \\ \text{---} \end{array} = \begin{array}{c} \text{---} \\ \boxed{f} \\ \text{---} \\ \boxed{h} \\ \text{---} \end{array} \quad \begin{array}{c} \text{---} \\ \boxed{g} \\ \text{---} \\ \boxed{w} \\ \text{---} \end{array} = (f \circ h) \otimes_1 (g \circ w)$$

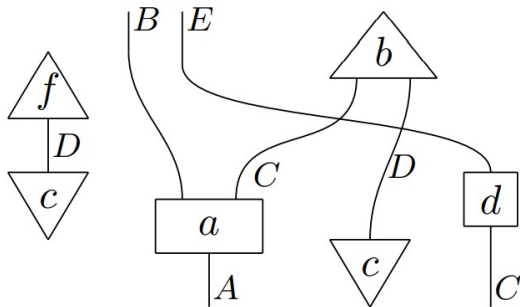
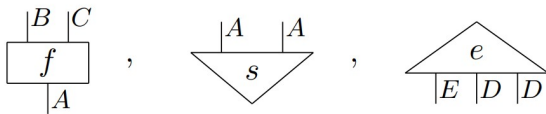
The diagram illustrates the equality of two tensor products of functions. On the left, the expression $(f \otimes_1 g) \circ (h \otimes_1 w)$ is shown. This is represented by two separate vertical stacks of boxes. The first stack has a box labeled f on top and a box labeled h on the bottom, connected by a vertical line. The second stack has a box labeled g on top and a box labeled w on the bottom, also connected by a vertical line. These two stacks are separated by an equals sign. To the right of this equals sign is another set of two vertical stacks of boxes. The first stack has a box labeled f on top and a box labeled h on the bottom, connected by a vertical line. The second stack has a box labeled g on top and a box labeled w on the bottom, also connected by a vertical line. These two stacks are separated by an equals sign. To the right of this equals sign is the expression $(f \circ h) \otimes_1 (g \circ w)$.

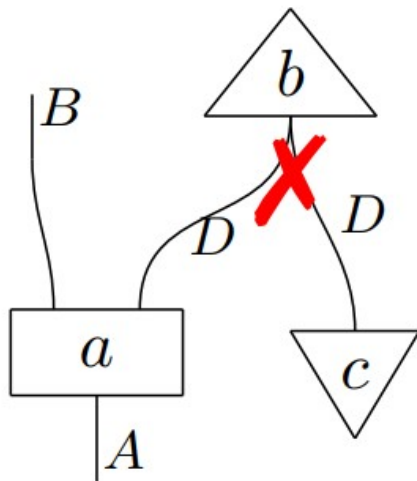


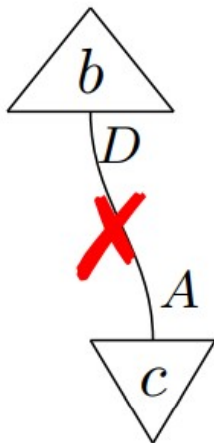
$$\text{SWAP}_{AB} := \begin{array}{c} \begin{array}{|c|c|} \hline B & A \\ \hline \end{array} \\ \text{---} \\ \begin{array}{|c|c|} \hline A & B \\ \hline \end{array} \end{array}$$

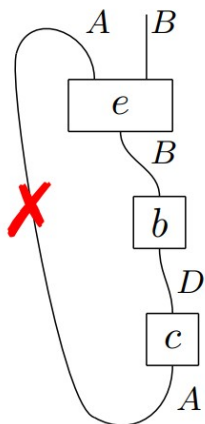












Thu 23rd July 2020 at 11.15

Ernesto F. Galvão (INL / UFF)

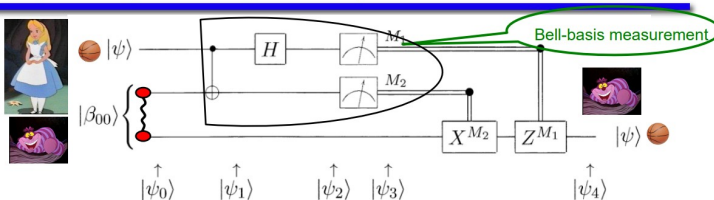
Quantum teleportation

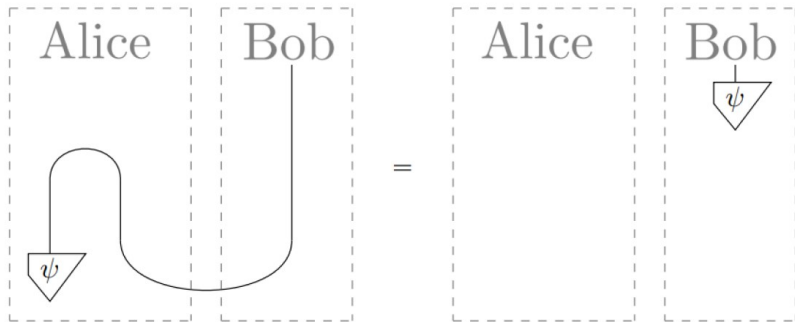
Abstract

Slides



Teletransporte quântico, passo a passo





Appendix



Theorem

Let $(H, \|\cdot\|)$ be a Hilbert space over $\mathbb{K} = \mathbb{R}, \mathbb{C}$. For any bounded $\phi : H \rightarrow \mathbb{K}$ ($\phi \in H^$) there is a unique vector $v_\phi \in H$ such that $\phi = \phi_{v_\phi}$ where $\phi_v := \langle v, \cdot \rangle$. For any bounded $\phi : H \rightarrow \mathbb{K}$, $\|v_\phi\| = \|\phi\|_\infty$. For all bounded $\phi, \phi' \in H^*$ and $\alpha \in \mathbb{K}$ it is true that $v_{\phi+\phi'} = v_\phi + v_{\phi'}$ and $v_{\alpha\phi} = \bar{\alpha}v_\phi$. In words, the mapping $\phi \mapsto v_\phi$ is antilinear.*

The proof goes as follows; for uniqueness, suppose that there are two vectors v_ϕ^1, v_ϕ^2 such that $\phi = \phi_{v_\phi^1} = \phi_{v_\phi^2}$. Then $\phi_{v_\phi^1} = \langle v_\phi^1, \cdot \rangle = \langle v_\phi^2, \cdot \rangle$ which happens iff $\langle v_\phi^1, v \rangle = \langle v_\phi^2, v \rangle$ for all $v \in H$. Let $v = v_\phi^1 - v_\phi^2$ we get $\langle v_\phi^1, v_\phi^1 - v_\phi^2 \rangle - \langle v_\phi^2, v_\phi^1 - v_\phi^2 \rangle = 0$ which implies that $v_\phi^1 - v_\phi^2 = 0$, hence uniqueness.



From the sesquilinearity of the inner product we have that $(\phi + \phi')(v) = \langle v_\phi + v_{\phi'}, v \rangle$ for all v therefore $v_{\phi+\phi'} = v_\phi + v_{\phi'}$ and $\phi(\alpha v) = \langle v_\phi, \alpha v \rangle = \overline{\alpha} \langle v_\phi, v \rangle$ hence $v_{\alpha\phi} = \overline{\alpha} v_\phi$.

Lastly, we show the relationship between the two norms. By definition we have

$\|\phi\|_\infty = \sup_{v \in H, \|v\|=1} |\phi(v)| = \sup_{v \in H, \|v\|=1} |\langle v_\phi, v \rangle| \leq \|v_\phi\|$ where in the last equality we use the Cauchy-Schwarz inequality. Therefore $\|\phi\|_\infty \leq \|v_\phi\|$. But we also have that,

$$\|v_\phi\| = \frac{|\langle v_\phi, v_\phi \rangle|}{\|v_\phi\|} = \left\langle v_\phi, \frac{v_\phi}{\|v_\phi\|} \right\rangle \leq \sup_{v \in H, \|v\|=1} |\langle v_\phi, v \rangle| = \|\phi\|_\infty$$



This result holds in general for any C^* -algebra \mathfrak{U} and the positive elements thereof \mathfrak{U}^+ . Therefore, one way to prove would be to use the theorem that for any polynomial p , the spectrum of operators A of the algebra satisfy that $\text{sp}(p(A)) = p(\text{sp}(A))$. We only need to choose the correct polynomial $p(x) := tx$ and the result follows for any real t .

Another way of showing this without using known results is as follows; considering only $t \geq 0$ and $t \in \mathbb{R}$. Let $t \neq 0$ at first.
 $z \in t\text{sp}(A) \iff \exists \lambda_z \in \mathbb{C}$ with $z = t\lambda_z$ such that $A - \mathbb{1}\lambda_z$ does not have inverse. This is true if, and only if $tA - \mathbb{1}t\lambda_z = tA - \mathbb{1}z$ does not have inverse, hence $t\text{sp}(A) = \text{sp}(tA)$ for $t \neq 0$. The case $t = 0$ follows from $\text{sp}(tA) = \{0\} = t\text{sp}(A)$.



We can state this for the more generic case of C^* -algebras, which have sets of matrices/hermitian operators as examples.

Lemma (generating)

Let \mathcal{A} be a C^ -algebra and $\mathcal{A}^{\mathbb{R}}$ the self-adjoint section. For any $A \in \mathcal{A}^{\mathbb{R}}$ there is A^+, A^- positive such that $A = A^+ - A^-$ and $\|A^+\|, \|A^-\| \leq \|A\|$.*

Proof.

The case $A = 0$ is true trivially. Let $A \neq 0$ then

$$A = \frac{\|A\|}{4} \left(\frac{A}{\|A\|} + \mathbb{1} \right)^2 - \frac{\|A\|}{4} \left(\frac{A}{\|A\|} - \mathbb{1} \right)^2 = p_1(A) - p_2(A).$$
 Where p_1, p_2 are two polynomials positive for every A , defining new self-adjoint operators $p_1(A), p_2(A)$ that have smaller norm and that must have positive spectrum, therefore they are positive.



$$\begin{pmatrix} a_0 + a_2 - \alpha & a_1 \\ a_1 & a_0 - a_2 - \alpha \end{pmatrix} \implies (a_0 + a_2 - \alpha)(a_0 - a_2 - \alpha) - a_1^2 = 0$$

$$a_0^2 - a_0 a_2 - \alpha a_0 + a_2 a_0 - a_2^2 - a_2 \alpha - a_0 \alpha + a_2 \alpha + \alpha^2 = a_1^2$$

$$\alpha(\alpha - 2a_0) = a_1^2 + a_2^2 - a_0^2 \implies a_1^2 + a_2^2 - a_0^2 \leq 0$$

The last implication is because for any positive real matrix σ we have $\text{Tr}(\sigma) = a_0$, since the trace is the sum of eigenvalues, any particular eigenvalue must be smaller than a_0 , hence $\alpha - 2a_0$ is always negative. The conclusion is then that for α to be nonnegative the rhs of the equation must be nonnegative.



First we show that if the probability of preparing a state is independent then the deterministic effect is unique. The probability of some state ω being prepared is $p(\omega|M) = \sum_{e \in M} e(\omega)$ in a given measurement M . Since the effects normalizing u_1, u_2 satisfy $u_1 = \sum_{e \in M} e$ and $u_2 = \sum_{e' \in M'} e'$ for two measurement procedures we have that $u_1(\omega) = p(\omega|M) = p(\omega|M') = u_2(\omega)$, because we assumed that the probability of being prepared was independent of future measurements to be applied, hence $p(\omega|M) = p(\omega) = p(\omega|M')$.



Second we show that if the deterministic effect is unique we get that the probability of preparing the state must be independent of future measurements. Now this direction is basically trivial from the above consideration. Whenever the deterministic effect is unique, given any two measurement procedures

$M = \{e\}, M' = \{e'\}$ we have that $u = \sum_{e \in M} e = \sum_{e' \in M'} e'$.

Therefore the probability of preparing the state ω will be equal in both cases since

$$p(\omega|M) = \sum_{e \in M} e(\omega) = u(\omega) = \sum_{e' \in M'} e'(\omega) = p(\omega|M').$$



Definition (Category)

Let C_0 and C_1 be two classes. The class C_0 is named the class of objects in the category and the class C_1 is named the class of arrows (or morphisms) in the category. There are two functions $\text{dom} : C_1 \rightarrow C_0$ and $\text{cod} : C_1 \rightarrow C_0$ such that if $f \in C_1$ we have that $\text{dom}(f), \text{cod}(f) \in C_0$ and we write $\text{dom}(f) \xrightarrow{f} \text{cod}(f)$.

Defining $H := \{(g, f) \in C_1 \times C_1 \mid \text{dom}(g) = \text{cod}(f)\}$ there is a function $\circ : H \rightarrow C_1$ and we write $g \circ f = \circ(g, f)$. There is a function $1 : C_0 \rightarrow C_1$. Then a category is a 6-tuple

$C = (C_0, C_1, \text{dom}, \text{cod}, \circ, 1)$ such that these structures satisfy the following demands:



- (a) If g and f are arrows, then $g \circ f$ is also an arrow. Therefore the composition operation \circ is closed within the category. This means that, as in the definition of the 6-tuple we have that $\text{cod}(g) = \text{dom}(f)$ implies that $\text{dom}(g \circ f) = \text{dom}(f)$ and $\text{cod}(g \circ f) = \text{cod}(g)$.
- (b) The map \circ is associative.
- (c) We have that $\text{dom}(1(A)) = \text{cod}(1(A))$ for every $A \in \mathbf{C}_0$ and $\forall f, g \in \mathbf{C}_1$ with $A \xrightarrow{f} \text{cod}(f)$ and $\text{dom}(g) \xrightarrow{g} A$ for any A it is true that 1_A does not changes the arrows, $f \circ 1_A = f$ and $1_A \circ g = g$.
- (d) **Lemma:** The map 1_A is unique for every $A \in \mathbf{C}_0$.



Definition

A strict monoidal category \mathbf{C} is a category such that there are two maps $\otimes_0 : \mathbf{C}_0 \times \mathbf{C}_0 \rightarrow \mathbf{C}_0$ and $\otimes_1 : \mathbf{C}_1 \times \mathbf{C}_1 \rightarrow \mathbf{C}_1$, and a unit object $I \in \mathbf{C}_0$ such that the following holds,

- (a) $\forall A, B, C \in \mathbf{C}_0$ we have that $(A \otimes_0 B) \otimes_0 C = A \otimes_0 (B \otimes_0 C)$.
- (b) $\forall A \in \mathbf{C}_0$ we have that $A \otimes_0 I = I \otimes_0 A = A$.
- (c) There is an interaction rule between \circ and \otimes_1 . They satisfy the so-called interchange law, so $\forall f_1, f_2, g_1, g_2 \in \mathbf{C}_1$ we have that $(g_1 \otimes_1 g_2) \circ (f_1 \otimes_1 f_2) = (g_1 \circ f_1) \otimes_1 (g_2 \circ f_2)$



Theorem: Every monoidal category is equivalent to a strict monoidal category. In particular, every monoidal category of sets has a strict monoidal structure, so there is no need to change the category structure.



A symmetric monoidal category (SMC) is a monoidal category with a swap morphism $\sigma_{A,B} : A \otimes B \rightarrow B \otimes A$ that is defined for any object satisfying:

- (a) Swapping twice is the identity, $\sigma_{A,B} \circ \sigma_{B,A} = 1_{A \otimes B}$.
- (b) Parallel composition followed by swapping is equal to swapping in the respectively type-objects and then composing in parallel with reversed order,
$$\sigma_{A,B} \circ (g \otimes f) = (f \otimes g) \circ \sigma_{C,D}.$$
- (c) Swapping the identity object with another object is the identity arrow, $\sigma_{A,I} = 1_A$.
- (d) This is somewhat a coherence relation between sapping and the identity arrows, $(1_B \otimes \sigma_{A,C}) \circ (\sigma_{A,B} \otimes 1_C) = \sigma_{A,B \otimes C}$.

